#### **2024 Workshop**

**Numerical Modeling of Earthquake Motions: Waves and Ruptures**

### **Finite-Difference Numerical Simulation of Seismic Waves Propagation in Models with Complex Boundary Geometries**

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#### ➢ **Finite-difference method and curvilinear grids**

### ➢ **How to implement FDM on curvilinear grids**

#### ➢ **How to handle complex geometries**

### **Finite-difference method (FDM)**

# $\rho \frac{\partial v_i}{\partial t} = \frac{\partial \sigma_{ij}}{\partial x_i} + f_i \ .$  $\left( \frac{\partial \sigma_{ij}}{\partial t} = \lambda \delta_{ij} \frac{\partial v_k}{\partial x_k} + \mu \! \left( \frac{\partial v_i}{\partial x_i} + \frac{\partial v_j}{\partial x_i} \right) \right)$

Governing equation **FD** discretization

$$
\dfrac{\dfrac{\partial v_i}{\partial t} \sim \dfrac{v_i^{n+1} - v_i^n}{\Delta t}}{\dfrac{\partial \sigma_{xy}}{\partial x} \sim \dfrac{\sigma_{xy} \Bigl( + \dfrac{\Delta x}{2} \Bigr) - \sigma_{xy} \Bigl( - \dfrac{\Delta x}{2} \Bigr)}{\Delta x}}{\dfrac{\partial \sigma_{xy}}{\partial y} \sim \dfrac{\sigma_{xy} \Bigl( + \dfrac{\Delta y}{2} \Bigr) - \sigma_{xy} \Bigl( - \dfrac{\Delta y}{2} \Bigr)}{\Delta y}}
$$

 $\Box$  theory is straightforward  $\Box$  gridded complex velocity structures  $\Box$  computational efficiency  $\Box$  easy to be used



### **Limitation of early FDM: can't deal with topography**



Surface topography Using Cartesian grid to approximate topography: staircase shape



- Zahradník and Urban, 1984
- Frankel and Leith, 1992,
- Robertsson, 1996
- Opršal and Zahradník, 1999
- Hayashi et al., 2001
- Zeng et al., 2012

Very dense grid is required (~60 PPW for surface wave) (Bohlen and Saenger, 2006)



### **Limitation of early FDM: can't deal with topography**



Surface topography Using Cartesian grid with immersed boundary treatment





Lombard et al., 2008 Gao et al., 2015

set values above free surface according to analytical condition hard to code for 3D and stability problem

#### **Limitation of Cartesian grids, but not FDM**



### **FDM not limited to Cartesian grid, but structural grids**

unstructured grid and structural grid



(Fichtner, 2011)



Cartesian grid

curvilinear grid boundary-conforming grid



- □ handle boundary and internal complex interfaces
- element conforming to the interfaces
- professional CAD software
- $\square$  professional grid generation software

 $\Box$  very easy to generate (x0, dx, nx) **□** dense grid to represent surface topography

- can represent surface topography well
- $\blacksquare$  very easy

to professional grid

generation software

#### **FDM only requires the grid to conform to the free surface but not the internal interfaces**

#### effective medium parameterization of the internal interfaces

- ➢ grid averaging (Graves, 1996)
- ➢ volume averaging (Moczo et al., 2002)
- ➢ Orthorhombic medium representation (Moczo et al., 2002, 2014, 2019; Kristek et al., 2017)
- ➢ TTI equivalent medium parameterization (Jiang and Zhang, 2021; Koene et al., 2022)



long-wave equivalent theory (Backus,1962)

#### **PPW~5 using TTI effective media**





#### **Efficient method to evaluate effective medium in 2D and 3D**



use an auxiliary grid to identify which point should be calculated

use Marching Cube method to calculate volume and normal direction

Jiang, L., and W. Zhang, Efficient implementation of equivalent medium parameterization in finite-difference seismic wave simulation methods, GJI, minor revision

https://github.com/jianglq6/FD3D\_ModelPrepare

### **TOC**

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- ➢ **How to implement FDM on curvilinear grids**
	- ➢ **Governing equations**
	- ➢ **FD schemes**
	- ➢ **Free surface boundary condition**
- ➢ **How to handle complex geometries**

## **Governing equation in curvilinear grid**

- Momentum equation:  $\rho \frac{\partial \mathbf{v}}{\partial t} = \nabla \cdot \mathbf{\sigma} + \mathbf{f}$ 
	- curvilinear grid represents a general curvilinear coordinate



physical coordinate  $(x,z)$ 

computational/  $(\xi,\zeta)$ curvilinear coordinate

 $\square$  spatial derivatives are evaluated along the curvilinear grid lines

- Wavefield variables can be defined in two different ways
	- 1. direction of components along the general coordinates
	- 2. direction of components along the global Cartesian coordinate

Governing equation:  $ho \frac{\partial v}{\partial t} = \nabla \cdot \sigma + f$ 

1. variables are defined in the general coordinates (Tensorial formulation )



Governing equation:  $ho \frac{\partial \mathbf{v}}{\partial t} = \nabla \cdot \mathbf{\sigma} + \mathbf{f}$ 

2. variables are defined in the global Cartesian coordinates



curvilinear coordinate

$$
\rho v_{x,t}\!=\!\xi_{,x}\sigma_{xx,\xi}\!+\!\zeta_{,x}\sigma_{xx,\zeta}\!+\!\xi_{,z}\sigma_{xz,\xi}\!+\!\zeta_{,z}\sigma_{xz,\zeta}
$$

#### Cartesian coordinate **Curvilinear** Curvilinear coordinate

$$
\rho v_{x,t} = \sigma_{xx,x} + \sigma_{xz,z} + f_x,
$$
\n
$$
\rho v_{x,t} = \sigma_{xx,x} + \sigma_{xz,z} + f_x,
$$
\n
$$
\rho v_{z,t} = \xi_{,x} \sigma_{xx,\xi} + \xi_{,z} \sigma_{zz,\xi} + \zeta_{,x} \sigma_{xx,\zeta} + \zeta_{,z} \sigma_{zz,\zeta},
$$
\n
$$
\rho v_{z,t} = \xi_{,x} \sigma_{xz,\xi} + \xi_{,z} \sigma_{zz,\xi} + \zeta_{,x} \sigma_{xz,\zeta} + \zeta_{,z} \sigma_{zz,\zeta},
$$
\n
$$
\sigma_{xx,t} = (\lambda + 2\mu)v_{x,x} + \lambda v_{z,z},
$$
\n
$$
\sigma_{zz,t} = \lambda v_{x,x} + (\lambda + 2\mu)v_{z,z},
$$
\n
$$
\sigma_{zz,t} = \mu(v_{x,z} + v_{z,x}),
$$
\n
$$
\sigma_{zz,t} = \mu(\zeta_{,x}v_{x,\xi} + \zeta_{,x}v_{x,\zeta} + \zeta_{,z}v_{z,\zeta},
$$
\n
$$
\sigma_{xz,t} = \mu(\zeta_{,z}v_{x,\xi} + \zeta_{,x}v_{z,\xi} + \zeta_{,z}v_{x,\zeta} + \zeta_{,x}v_{z,\zeta}).
$$

$$
\mathbf{W}_{,t} = \mathbf{A}\mathbf{W}_{,x} + \mathbf{B}\mathbf{W}_{,z} \qquad \qquad \mathbf{W}_{,t} = \mathbf{A}\mathbf{W}_{,\xi} + \mathbf{B}\mathbf{W}_{,\zeta}.
$$

the structure of the equations is similar, with additional differential terms

$$
\rho v_{x,t}\!=\!\xi_{,x}\sigma_{xx,\xi}\!+\!\zeta_{,x}\sigma_{xx,\zeta}\!+\!\xi_{,z}\sigma_{xz,\xi}\!+\!\zeta_{,z}\sigma_{xz,\zeta}
$$

staggered-grid schemes: variables are defined at staggered positions

 $\sigma_{xx}, \sigma_{zz}, c_{ii}$  $\sigma_{xz}, c_{ij}$  $v_x, \rho$  $\bigtriangledown$  $v_z, \rho$ SSG

#### interpolation

Lebedev grid staggered grid
staggered grid
staggered grid
staggered grid
staggered grid
staggered
stagger



calculate by another staggered grid

Igel et al. (1995) Lisitsa and Vishnevskiy (2010)

#### supporting operator mimetic scheme



pair of gradient and divergence operators

Ely et al. (2006); Shragge and Tapley (2017)

$$
\rho v_{x,t}\!=\!\xi_{,x}\sigma_{xx,\xi}\!+\!\zeta_{,x}\sigma_{xx,\zeta}\!+\!\xi_{,z}\sigma_{xz,\xi}\!+\!\zeta_{,z}\sigma_{xz,\zeta}
$$

collocated-grid schemes: variables are defined at same positions



$$
\rho v_{x,t}\!=\!\xi_{,x}\sigma_{xx,\xi}\!+\!\zeta_{,x}\sigma_{xx,\zeta}\!+\!\xi_{,z}\sigma_{xz,\xi}\!+\!\zeta_{,z}\sigma_{xz,\zeta}
$$

collocated-grid schemes: variables are defined at same positions



$$
\rho v_{x,t}\!=\!\xi_{,x}\sigma_{xx,\xi}\!+\!\zeta_{,x}\sigma_{xx,\zeta}\!+\!\xi_{,z}\sigma_{xz,\xi}\!+\!\zeta_{,z}\sigma_{xz,\zeta}
$$

collocated-grid schemes: variables are defined at same positions



$$
\rho v_{x,t}\!=\!\xi_{,x}\sigma_{xx,\xi}\!+\!\zeta_{,x}\sigma_{xx,\zeta}\!+\!\xi_{,z}\sigma_{xz,\xi}\!+\!\zeta_{,z}\sigma_{xz,\zeta}
$$

collocated-grid schemes: variables are defined at same positions



3 solutions of odd-even decoupling on collocated-grid FD scheme:

(1) centered difference + explicit filtering (2) centered difference +inherent dissipation

Lax-wendroff scheme

$$
\partial_t u + \partial_x f(u) = 0
$$
\n
$$
\begin{aligned}\nu(x_i, t^{n+1}) &= u(x_i, t^n) + \\
-\Delta t a \partial_x^{(1)} u(x_i, t^n) + \frac{\Delta t^2}{2} a^2 \partial_x^{(2)} u(x_i, t^n) + O(\Delta t^2) \\
&\uparrow\n\end{aligned}
$$
\ndissipation term

(Bogey & Bailly 2006) https://encyclopediaofmath.org/wiki/Lax-Wendroff\_method

$$
f(\boldsymbol{U}_i)^d = \boldsymbol{U}_i - \sigma d_0 \boldsymbol{U}_i - \sigma \sum_{m=1}^6 d_m (\boldsymbol{U}_{i+m} + \boldsymbol{U}_{i-m})
$$

 $a_m (U_{i+m} - U_{i-m})$  ,

 $\partial \bm{U}_i$ 

 $\partial \xi$ 

=

1

 $\sum$ 

6

 $m=1$ 

 $\Delta \xi$ 

filter to remove high wavenumber

3 solutions:

(3) MacCormack-type schemes



$$
D_x^F v_{x|i} = \frac{1}{\Delta x} (a_3 v_{x|i+3} + a_2 v_{x|i+2} + a_1 v_{x|i+1} + a_0 v_{x|i} + a_{-1} v_{x|i-1})
$$
biased forward operator  
\n
$$
D_x^B v_{x|i} = \frac{1}{\Delta x} (\n\begin{array}{cc}\n-a_{-1} v_{x|i+1} - a_0 v_{x|i} - a_1 v_{x|i-1} - a_2 v_{x|i-2} - a_3 v_{x|i-3}\n\end{array})
$$
biased backward operator  
\n
$$
D_x v_{x|i} = \frac{1}{2} (D_x^F v_{x|i} + D_x^B v_{x|i})
$$
if second order RK used  
\ncenter difference is recovered

#### asymmetric stencil has both dispersion and dissipation errors

(Bayliss et al., 1986; Zhang and Chen, 2006; Zhang et al., 2012)

### **Free surface boundary condition**

#### **characteristic boundary condition**

decomposition of the wave equation into incoming and outgoing wave modes at subdomain boundaries

$$
\begin{array}{ccc}\n & \sigma_{xz}^{(new)} = 0 & \sigma_{xz}^{(new)} = 0 \\
\hline\n\mathbf{W}_{,t} = \mathbf{A}\mathbf{W}_{,x} + \mathbf{B}\mathbf{W}_{,z} & \uparrow c_s & \uparrow c_p & \uparrow c_p & \downarrow c_p \\
\mathbf{B} = \mathbf{S}\mathbf{\Lambda}\mathbf{S} & \stackrel{\frown}{\longrightarrow} & R_1 = \sqrt{\rho\mu} \, v_x + \sigma_{xz} & \downarrow c_p & \downarrow c_x \\
 & R_2 = \sqrt{(\lambda + 2\mu)\rho} \, v_z + \sigma_{zz} & \downarrow c_x & \downarrow c_x \\
 & R_3 = (\lambda + 2\mu)\sigma_{xx} - \mu\sigma_{zz} & v_z^{(new)} = v_z - \frac{\sigma_{zz}}{\sqrt{(\lambda + 2\mu)\rho}} \\
\end{array}
$$

 $(n \text{ on})$ 

(Gottlieb, 1982; Bayliss et al., 1986; Carcione, 1991)

### **Free surface boundary condition**

#### **Traction imaging technique**

$$
\rho v_{x,t} = \xi_{,x} \sigma_{xx,\xi} + \xi_{,z} \sigma_{xz,\xi} + \zeta_{,x} \sigma_{xx,\zeta} + \zeta_{,z} \sigma_{xz,\zeta},
$$
  
\n
$$
\left[\begin{array}{c}\n\text{rewrite using conservative form of } \nabla \\
\varphi \frac{\partial v_x}{\partial t} = \frac{1}{J} \left\{ \frac{\partial}{\partial \xi} \left[ J \left( \frac{\partial \xi}{\partial x} \sigma_{xx} + \frac{\partial \xi}{\partial z} \sigma_{xz} \right) \right] \right. \\
\left. + \frac{\partial}{\partial \zeta} \left[ J \left( \frac{\partial \zeta}{\partial x} \sigma_{xx} + \frac{\partial \zeta}{\partial z} \sigma_{xz} \right) \right] \right\} \\
\approx T_x\n\end{array}\right]
$$



(Zhang and Chen, 2006; Zhang et al, 2012)

### **Smooth topography: vertical stretched grid**



SEM: black line

### **Topography+fault: general curvilinear grid**



## **Rough topography: general curvilinear grid**





For seismic wave numerical simulation, only free surface is fixed, shapes of the other boundaries are free to be chosen



#### oblique grid does not affect accuracy, but reduce the time step of numerical stability



#### **Global seismic waveform simulation (2D) with surface topography**



physical coordinates of the grid

 $x(\xi, \eta) = R(\xi, \eta) * \sin(\theta(\xi)),$ 

 $z(\xi, \eta) = R(\xi, \eta) * \cos(\theta(\xi)),$ 

θis the longitude, R is the radius, which can vary with longitude

➢variables are defined in global cartesian coordinate  $\triangleright$  gradients are calculated in curvilinear coordinate  $\triangleright$  input and output: polar coordinate

A special curvilinear grid



#### Free surface topography can be taken into account



#### **Do not define point at center multi-level discontinues-grid**





filter (Kristek et al., 2010; Zhang et al., 2013) high-order interpolation

### A simplified lunar model with high-speed anomalies and surface undulations





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### **Complex geometries: overset grid / multiblock grid**

#### Examples from computational Fluid Dynamics



(Bill Henshaw, 2011)

- ➢ Update wavefield on each block
- ➢ use high-order interpolation to exchange messages between blocks (set values at ghost points)

### **Overset grid for rough topography**

![](_page_35_Figure_1.jpeg)

➢ **reduce difficulty of grid generation:** fewer layers of curvilinear grid ➢ **improve computational efficiency:**

□ computing on Cartesian grids more efficient  $\square$  near orthogonal curvilinear grid: increase step size

Nan Zang, Wei Zhang, and Xiaofei Chen, (2021), An overset-grid finite-difference algorithm for simulating elastic wave propagation in media with complex free-surface topography, GEOPHYSICS, 86(4):T277-T292

![](_page_36_Figure_0.jpeg)

![](_page_36_Figure_1.jpeg)

6<sup>th</sup>-order Lagrange interpolation

$$
W_{target}(x_0, z_0) = \sum_{i=1}^{N=6} \sum_{j=1}^{N=6} L_{ij} W_{donor}(x_i, z_j)
$$

$$
L_{ij} = \prod_{l=1, l \neq i}^{N} \frac{x_0 - x_l}{x_i - x_l} \prod_{k=1, k \neq j}^{N} \frac{z_0 - z_k}{z_j - z_k}
$$

5 interior points + 3 ghost points overlaid

#### A tall hill model

![](_page_37_Figure_1.jpeg)

#### **Complex velocity model: Foothills**

![](_page_38_Figure_1.jpeg)

hlum

 $\mathsf 0$ 

Imm

 $\mathcal{R}$ 

 $Time(s)$ 

5

### **Conclusions and Further Works**

- Curvilinear grid FDM can simulate seismic wave propagation in models with complex boundary geometries
- Rough topography and whole Earth model with topography demonstrate the ability of the FDM to handle surface topography
- For very complex geometries, overset grid technique combining multiple cartesian and curvilinear grids is a standard and efficient solution for FDM
- Automatically generate 2D/3D curvilinear grid on the fly for large-scale simulation
- High-order free surface boundary condition implementation to reduce vertical PPW same to that of interior region

## **Thanks!**