



南京大學
NANJING UNIVERSITY

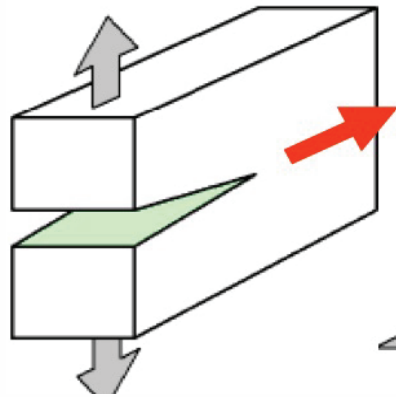
Viscoelastic ruptures unbounded by classical speed limits

Huihui Weng

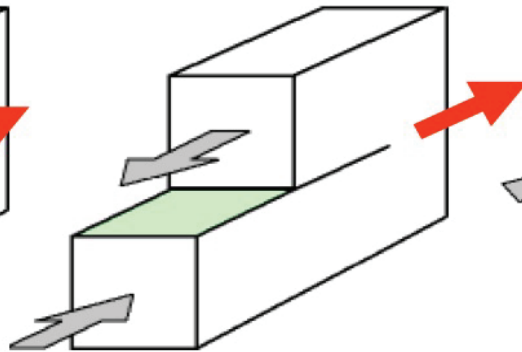
Monday, June 24, 2024
Smolenice Castle, Slovakia

Rupture modes

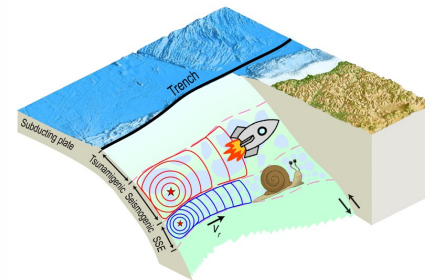
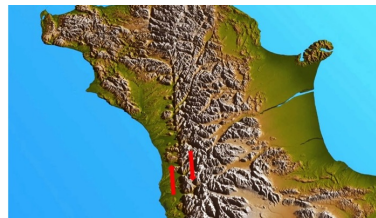
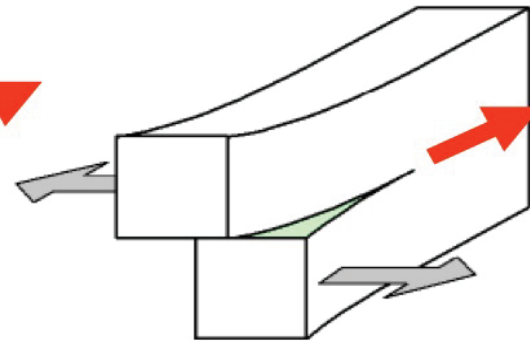
Mode I



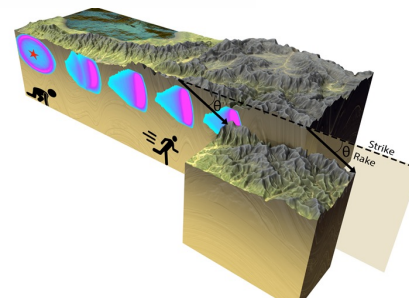
Mode II



Mode III

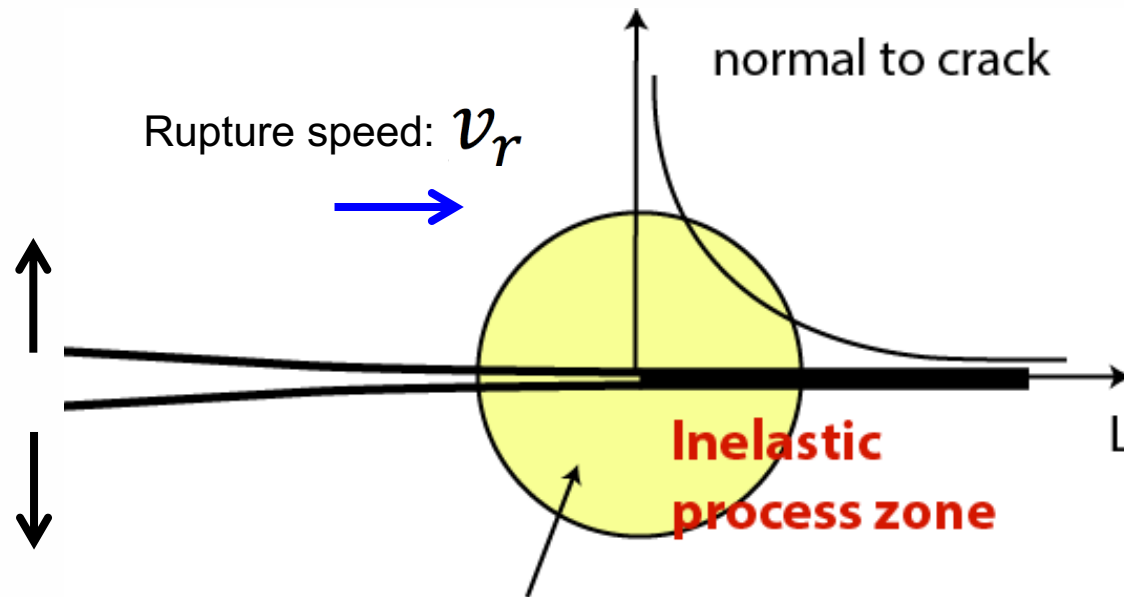
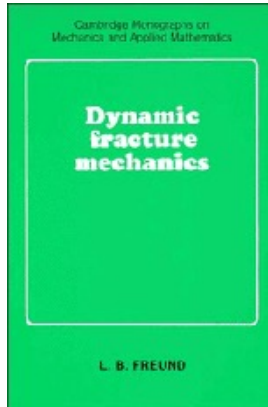


Mode -I + Mode II



Mode II + Mode III

2D Linear Elastic Fracture Mechanics



G : the elastic energy flowing into process zone

2D equation-of-motion: $G_c = G(v_r, L, \Delta\tau) \propto L$

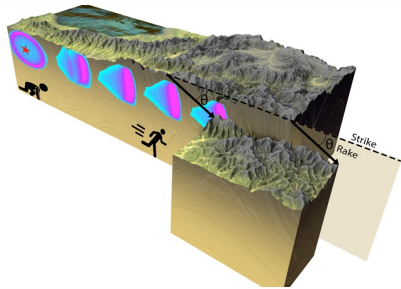
Fracture energy

Energy release rate

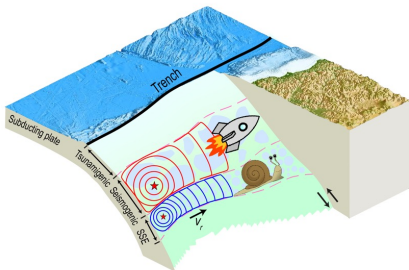
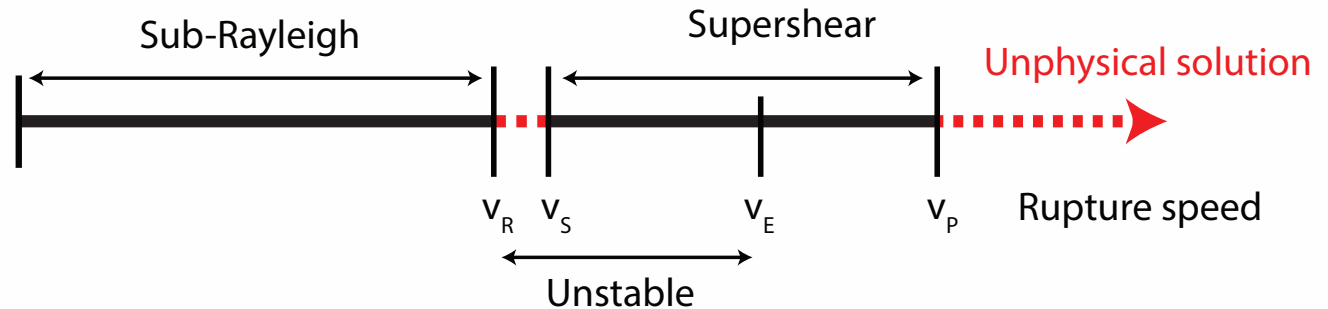
Classical speed limits



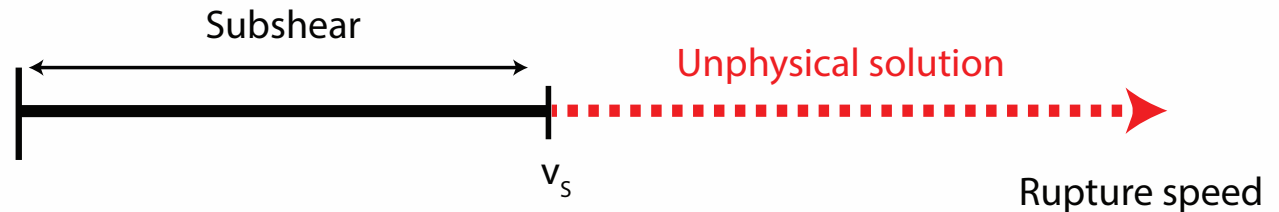
Modes I (dike)



Modes II (strike slip)

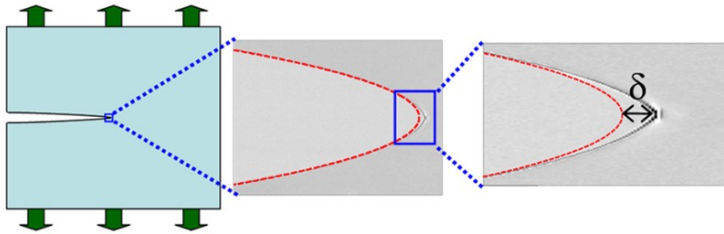


Modes III (dip slip)

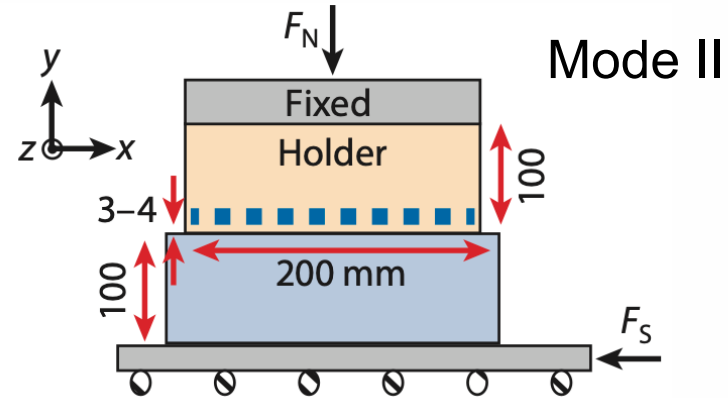
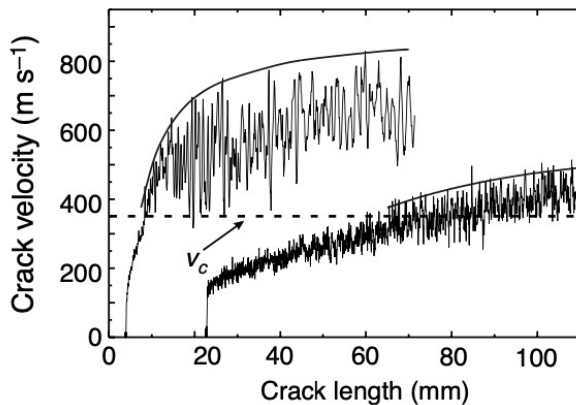


Validated speed limits in laboratory

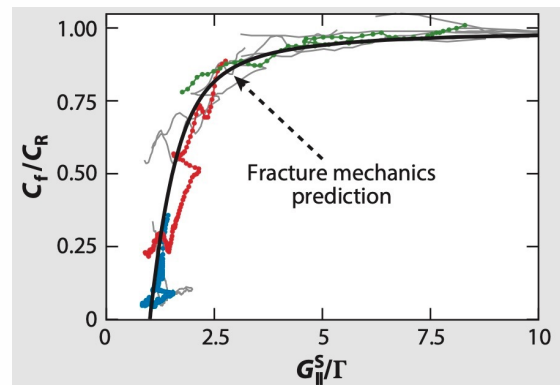
Mode I



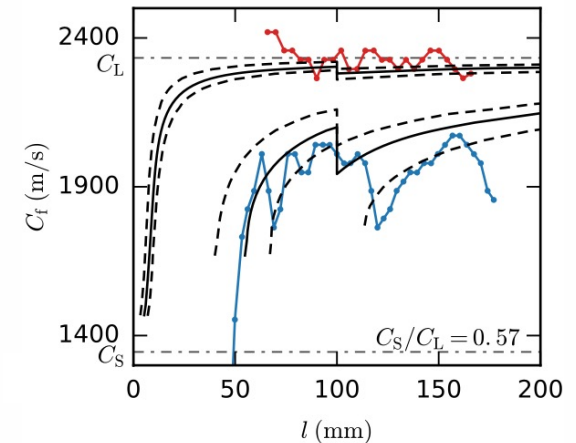
Subshear ($v < v_R$)



Subshear ($v < v_R$)



Supershear ($v < v_P$)

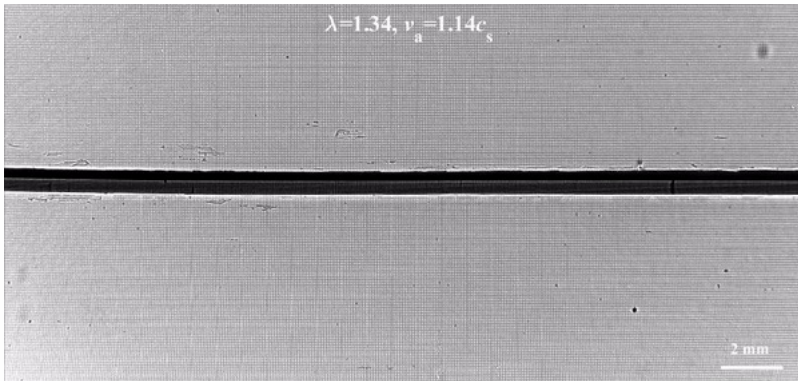


Beyond classical speed limits in 2D



Mach cones

Supershear crack in Mode I

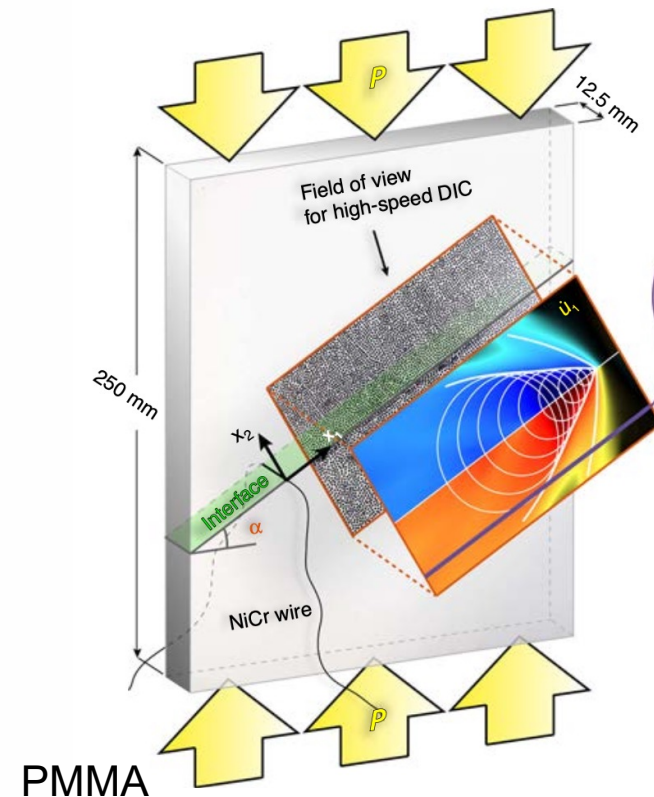


Hydrogel

Hyperelastic behavior

Wang et al, 2023

Supersonic crack in Mode II

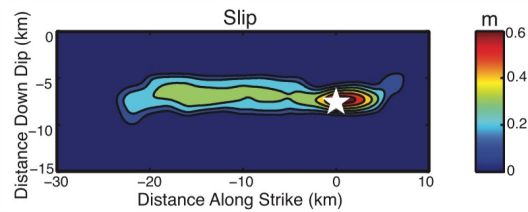


Viscoelastic behavior

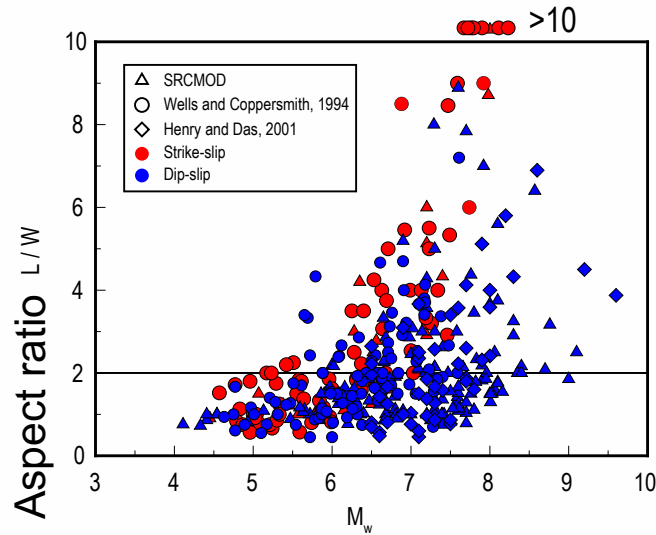
Gori et al, 2018

Finite rupture width in 3D

2004 Mw 6 Parkfield

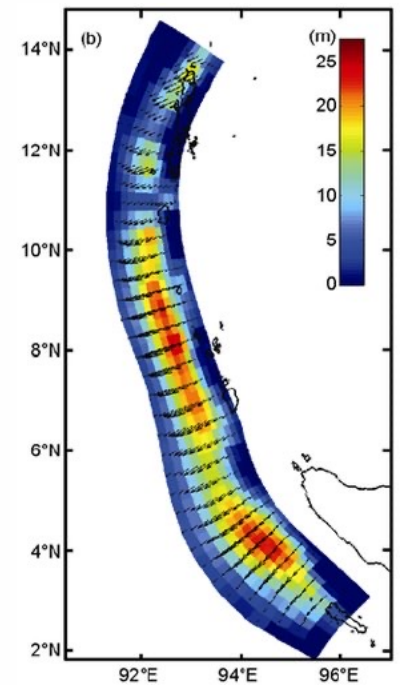


Ma et al 2008

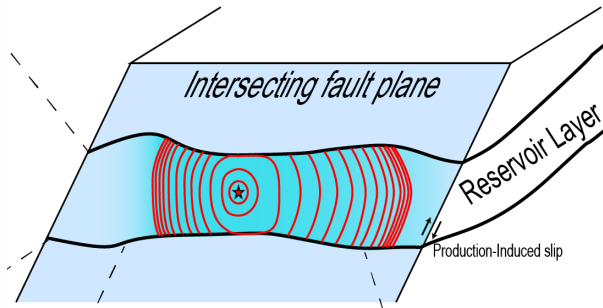


Weng and Yang, 2017

2004 Mw 9.3 Sumatra



Wang et al 2011



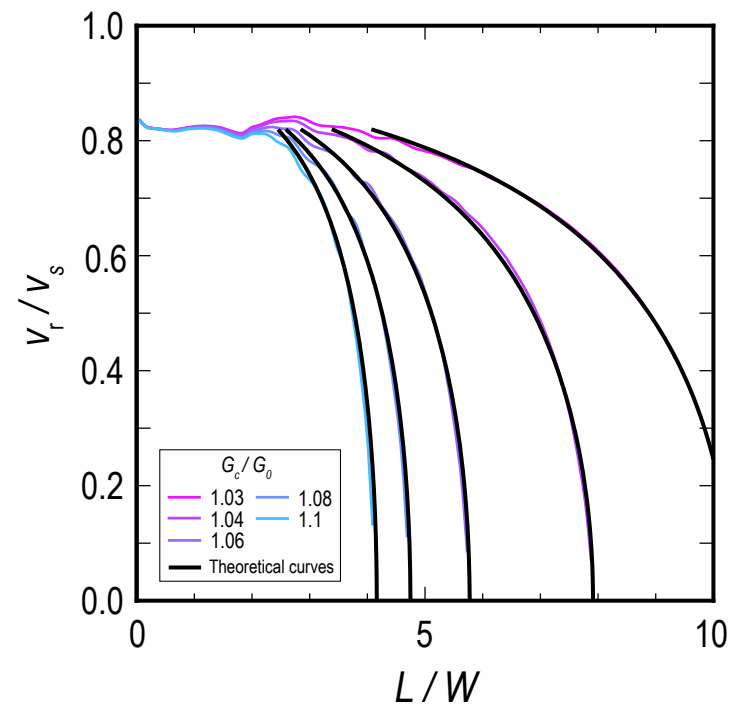
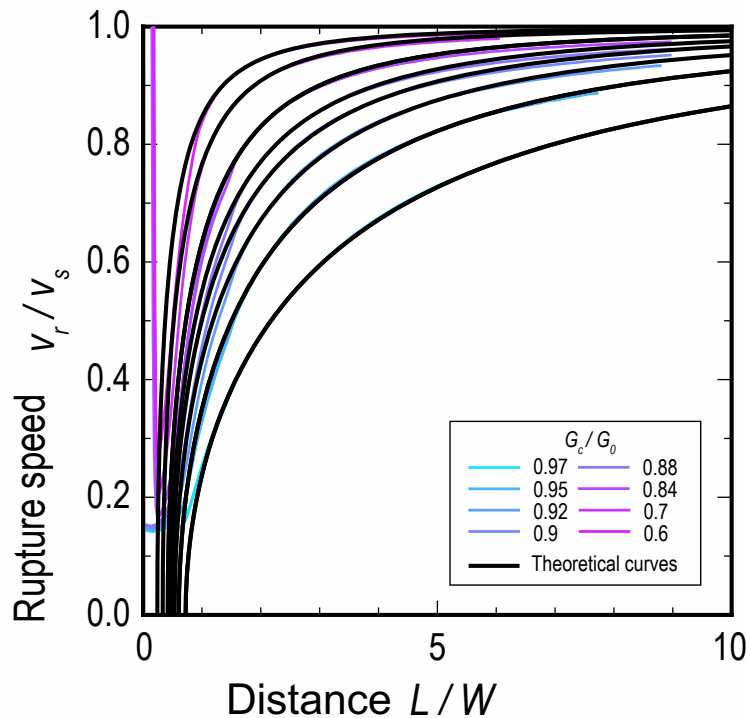
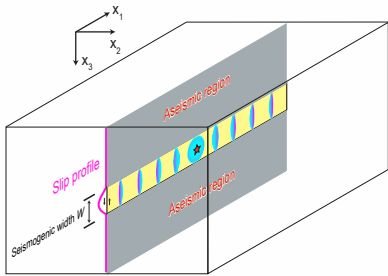
Induced seismicity

Extended 3D LEFM theory

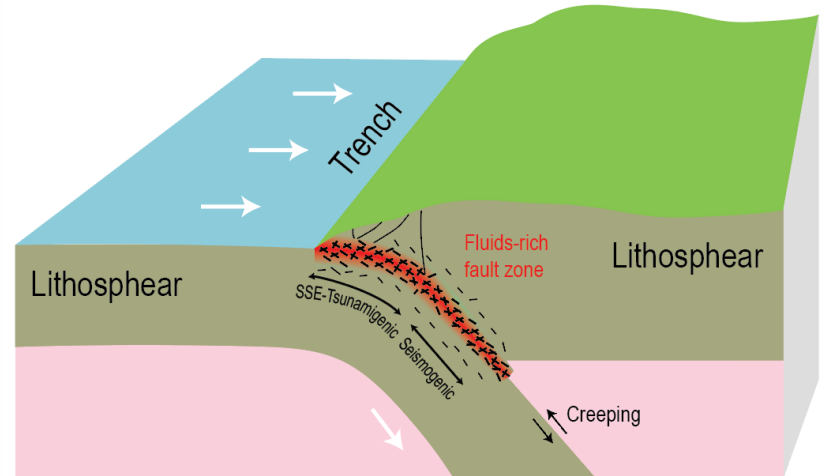
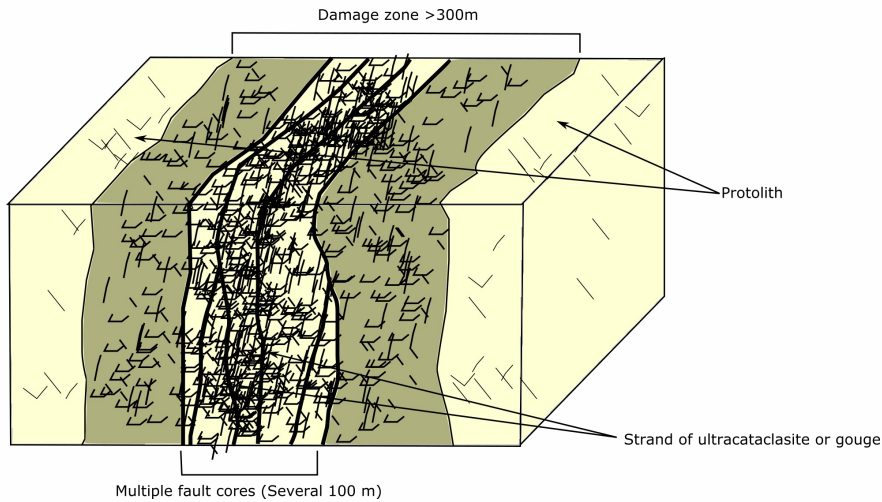
3D equation-of-motion:

$$F(G_c/G_0) = M(v_r) \cdot \dot{v}_r$$

"Force"
"Mass"
Acceleration



Highly damaged fault zone

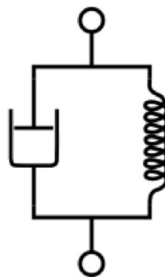


WordPress.com

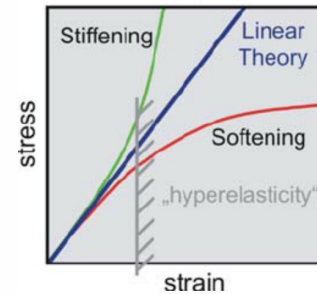
Elasticity?



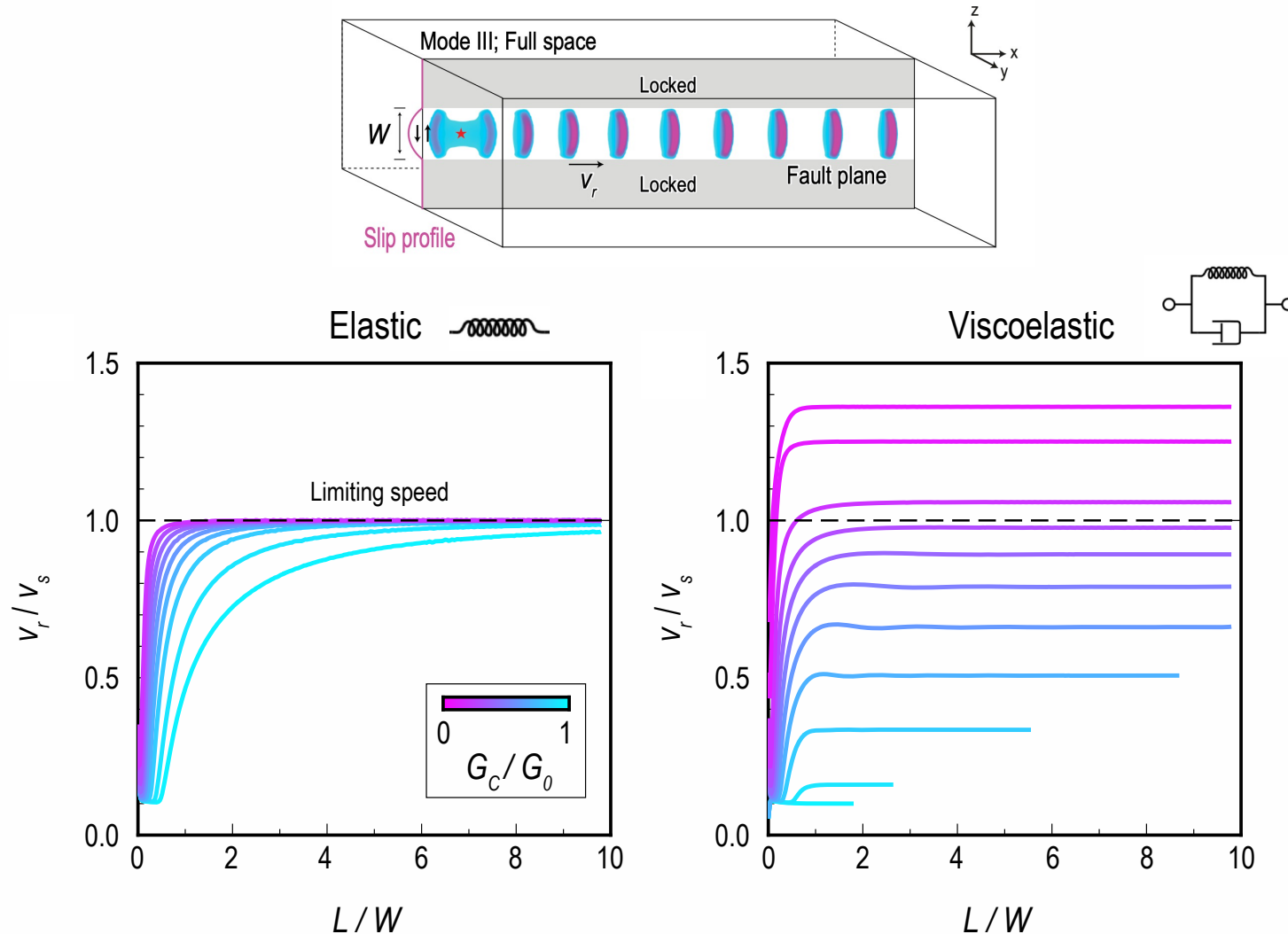
Viscoelasticity?



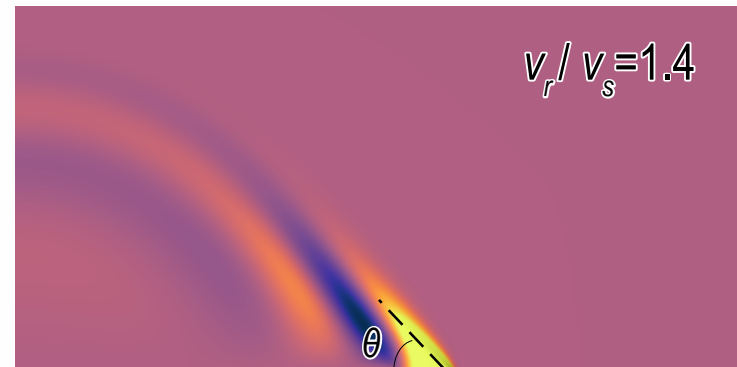
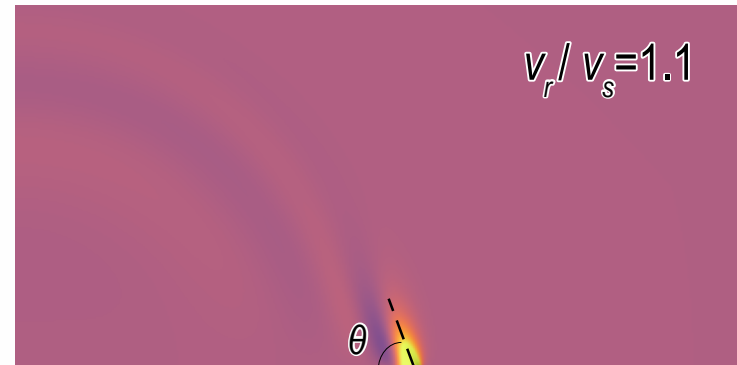
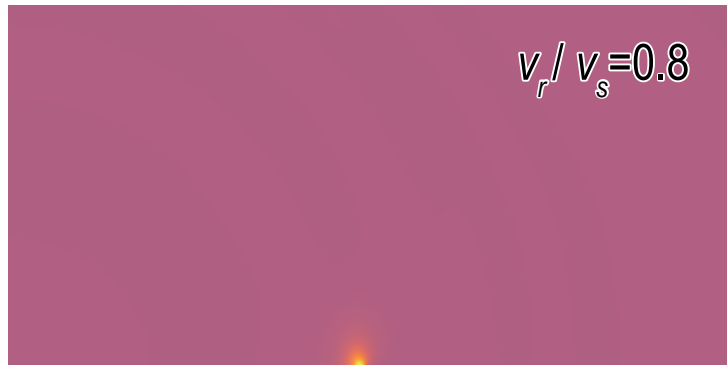
Hyperelasticity?



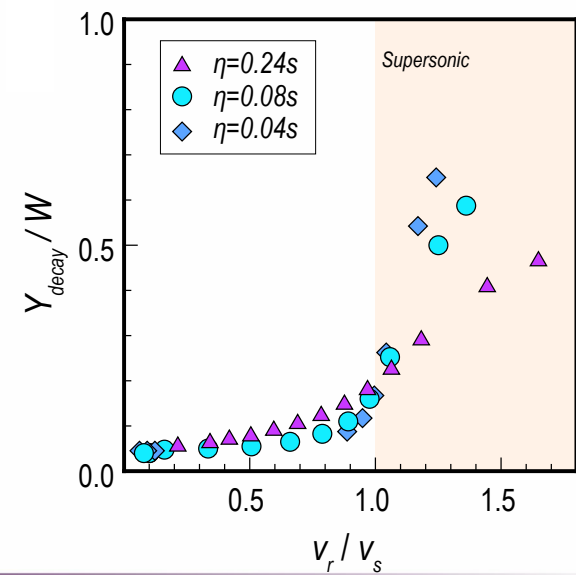
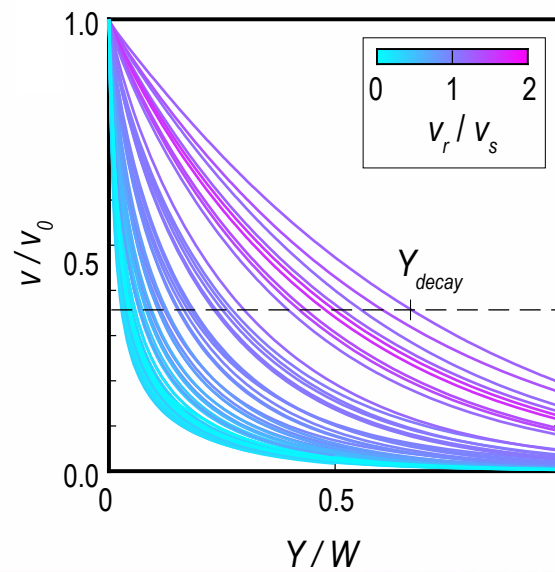
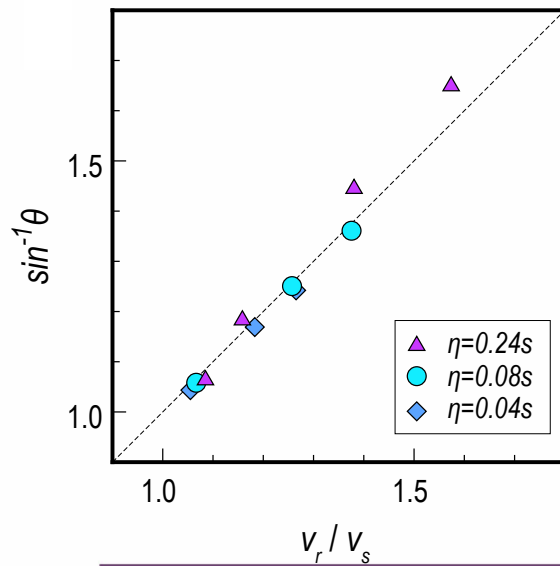
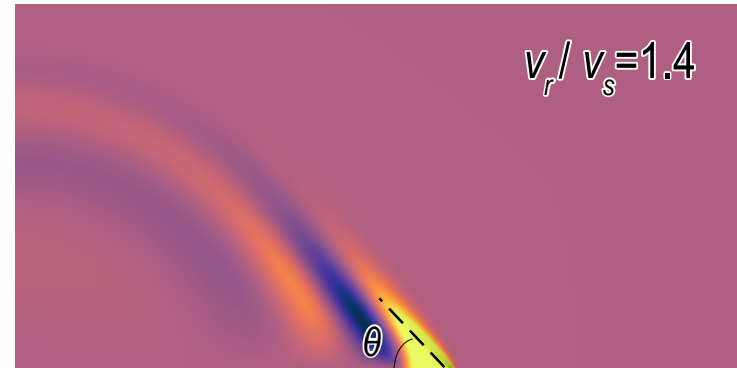
Unbounded ruptures in numerical simulations



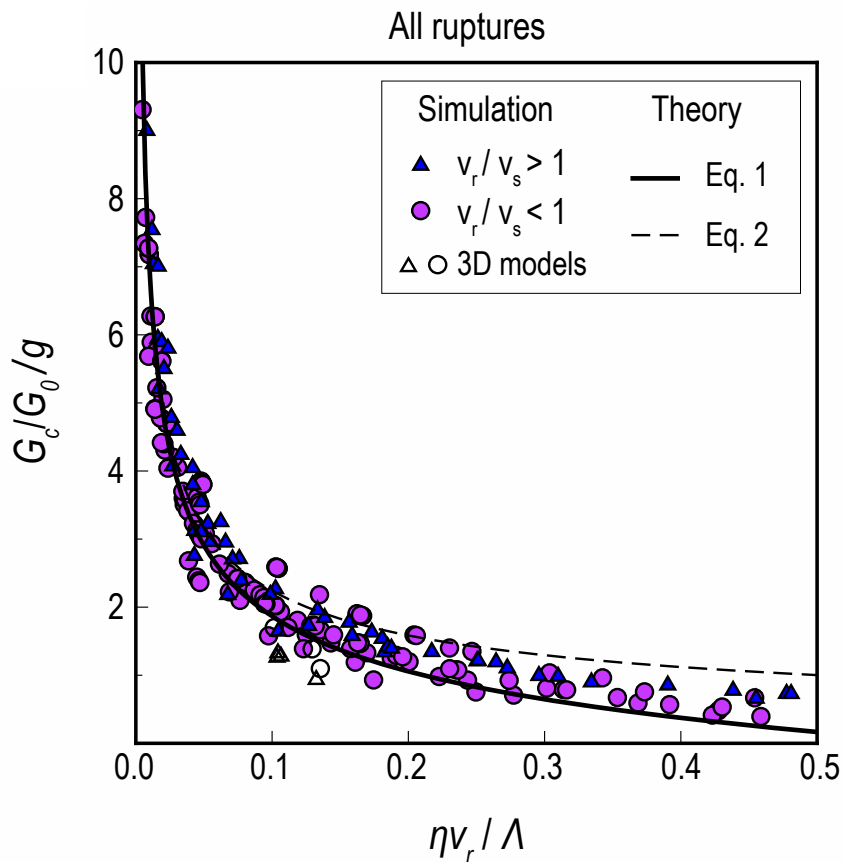
Mach fronts in mode III



Mach fronts in mode III



Theory for viscoelastic ruptures



Elastic model

$$F(G_c/G_0) = M(v_r) \cdot \dot{v}_r$$

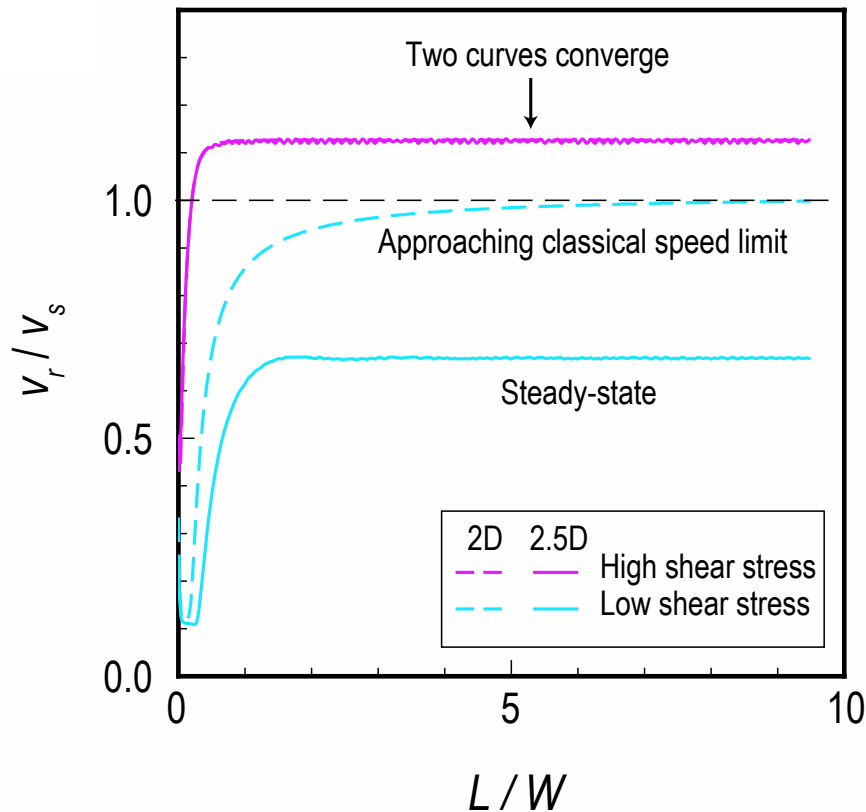
Extended

Viscoelastic model

$$\frac{G_c}{G_0} = \Theta\left(\frac{\eta v_r}{\Lambda}\right) \cdot g\left(\frac{\eta v_r}{W}\right)$$

Theory for viscoelastic ruptures

Also valid in 2D

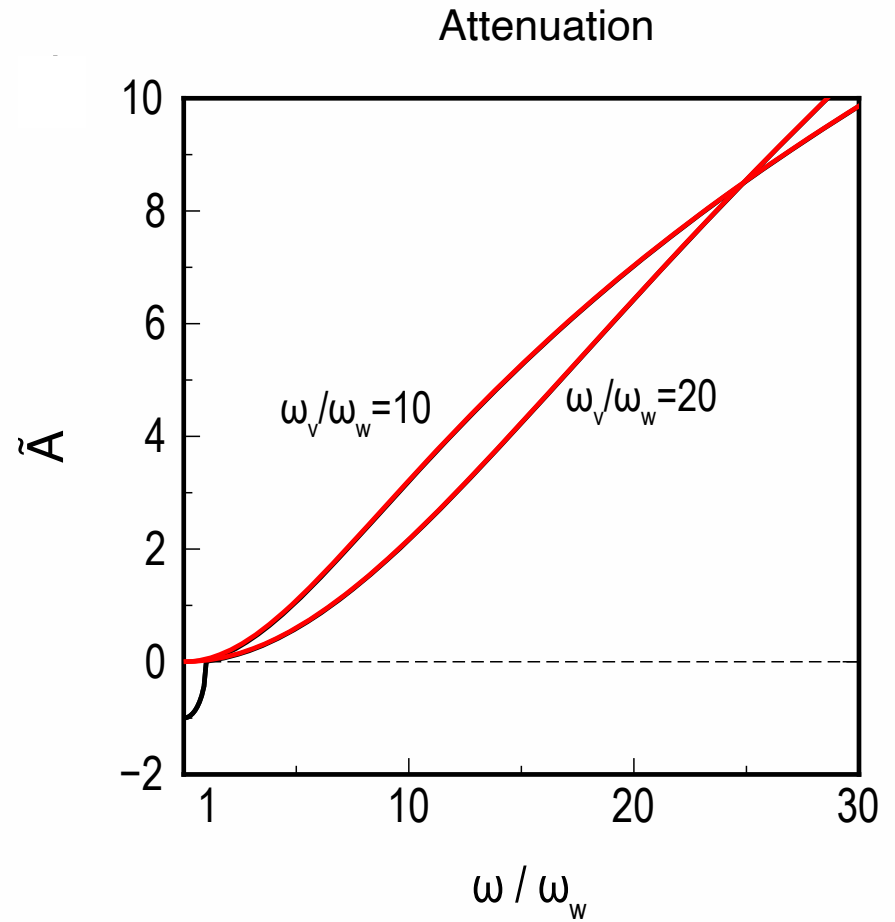
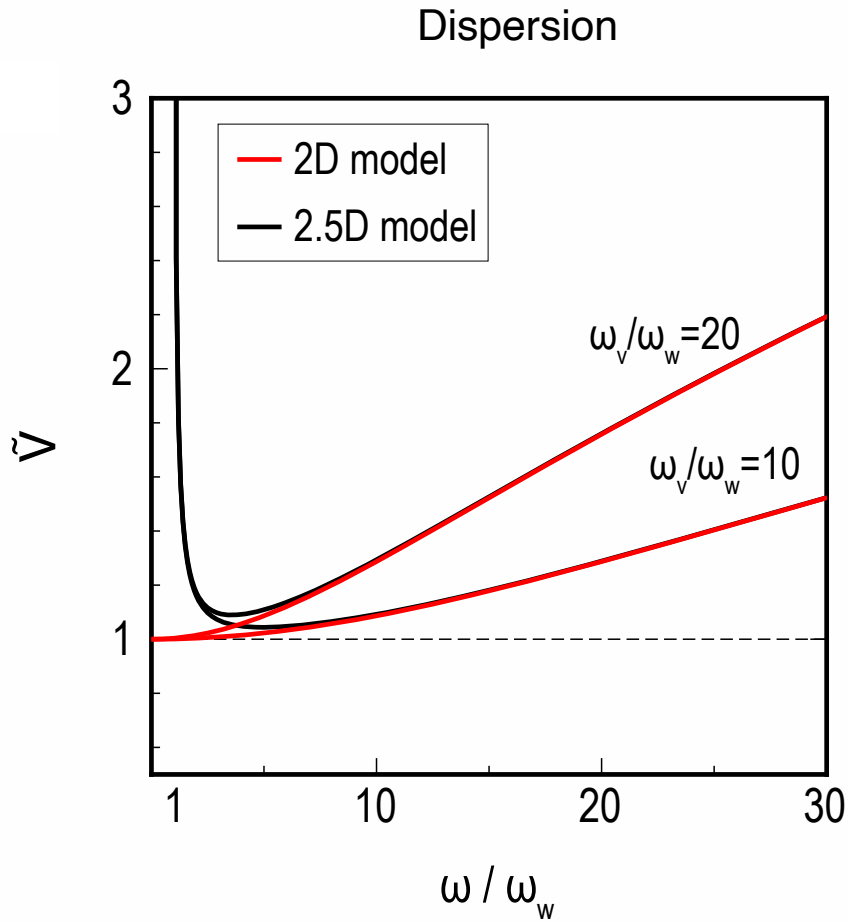


Asymptotically solution for supershear ruptures:

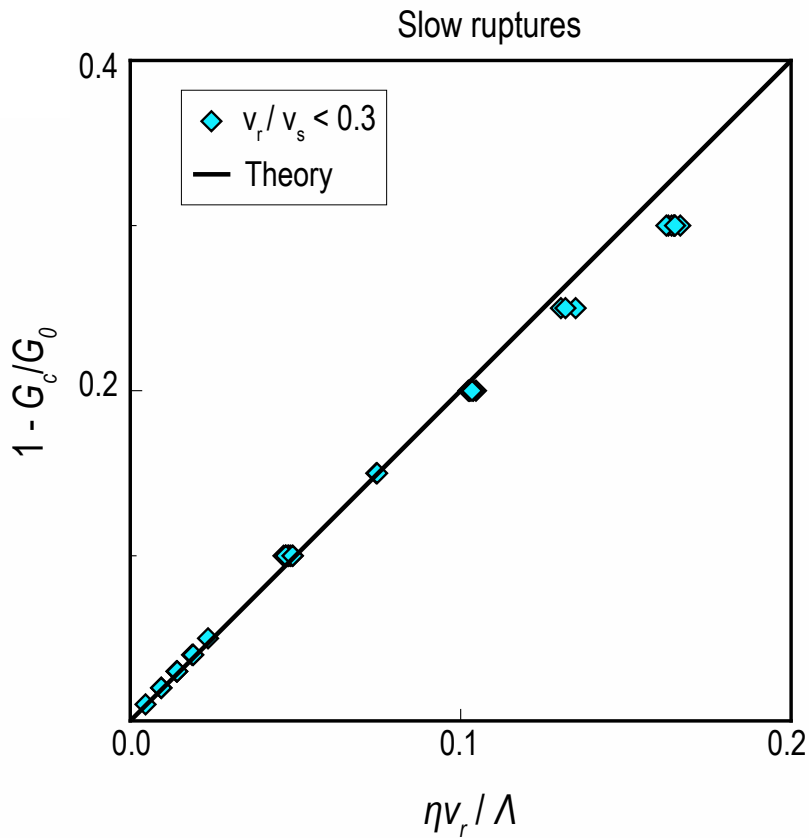
$$G_c \propto \frac{\Delta\tau^2 \sqrt{\eta v_r \Lambda}}{\hat{\alpha}_s \mu}$$

$$\hat{\alpha}_s = \sqrt{(v_r/v_s)^2 - 1}$$

Velocity dispersion and attenuation



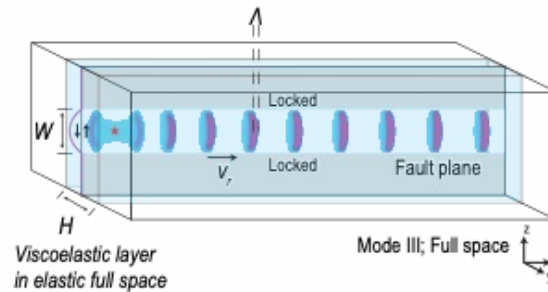
Theory for viscoelastic ruptures



Asymptotically solution for very slow ruptures:

$$G_0 = G_{equiv} \approx G_c \left(1 + 2 \frac{\eta v_r}{\Lambda} \right)$$

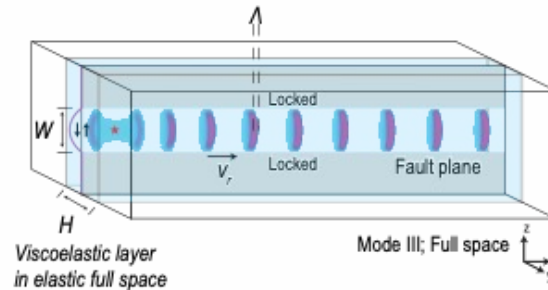
Finite thickness of viscoelastic layer



$H \rightarrow 0$ Purely elastic model

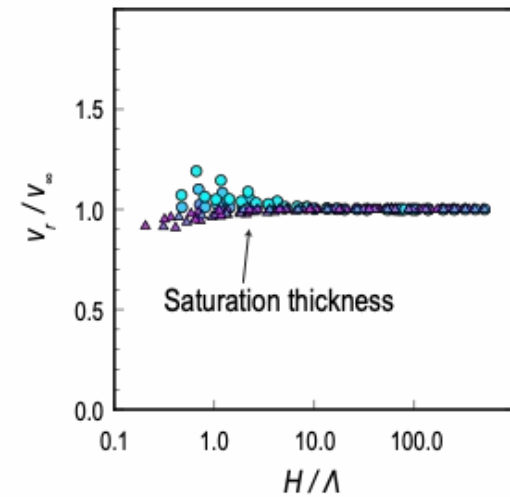
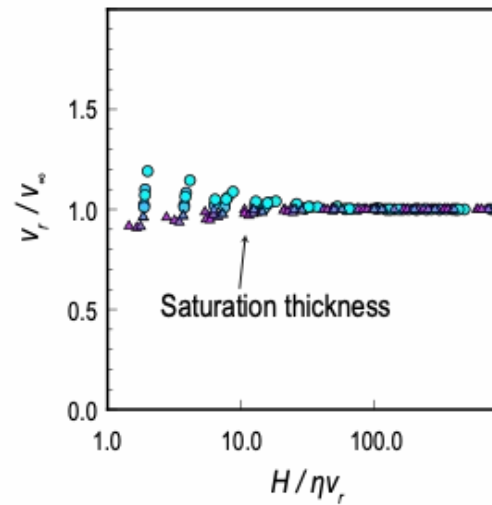
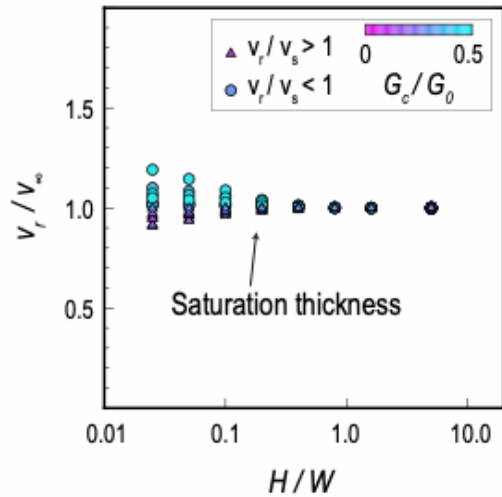
$H \rightarrow \infty$ Purely viscoelastic model

Finite thickness of viscoelastic layer



$H \rightarrow 0$ Purely elastic model

$H \rightarrow \infty$ Purely viscoelastic model



Take-home messages

- Viscoelastic ruptures can propagate at a continuum of terminal speeds not bounded by classical speed limits.
 - All simulated speeds are predicted by the new theory incorporating viscoelasticity.
 - Beyond classical speed limits, rupture dynamics are independent of any macroscopic length and is controlled only by local properties around the rupture tip.
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