

**New advances  
in implementing material heterogeneity  
in the finite-difference modelling**

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# Outline

- FDM
  - FDM – what is it ?
  - FDM among other numerical methods
  - Key conditions for competitiveness of FDM
- The problem: a discrete representation of heterogeneity of a medium in a heterogeneous FD scheme
  - Physics and mathematics of a discrete representation
  - Short excursion into a strange (though not unusual) and illuminating history
- The recent most advanced approaches
- Conclusions

## FDM

as applied to seismic wave propagation  
and earthquake ground motion

= a large diverse family

of computational schemes

based on

FD approximations

of the equation of motion and constitutive law

at space-time grid points

a state-of-the-art FD scheme

can significantly differ from some other FD scheme

in accuracy and computational efficiency

in a strongly heterogeneous medium

FDM, SEM and DGM

are the most important recent numerical-modelling methods

FDM is clearly dominant  
in the seismic prospecting

It could be also dominant

(due to accuracy\_and\_efficiency)

in investigation of earthquake ground motion

in local surface sedimentary structures

if

all the schemes being used

were at the state-of-the-art level

## Key conditions for competitiveness of FDM

an **explicit heterogeneous** FD scheme  
on a **uniform** spatial grid

(the latter does not contradict the use  
of an efficient discontinuous grid  
composed of several uniform grids)

**efficiency\_and\_accuracy**  
are determined by the  
**grid dispersion**  
and  
**discrete grid representation**  
of a material heterogeneity, mainly material interfaces

explicit

the field variable at a space-time grid point is calculated using an explicit FD formula that uses only values of the field variables at previous time levels

heterogeneous

one scheme is used for all interior grid points no matter what their positions are with respect to material interfaces

presence of interfaces is accounted for only by values of the effective material parameters assigned to grid positions

uniform-grid

for a chosen computational region and maximum frequency one grid can be used for arbitrary alterations of geometry/position of material interfaces

Obviously,  
a heterogeneous FD scheme  
should approximate  
an equation of motion and constitutive law  
valid and having the same form  
at any point of a medium

--

away of an interface or at an interface

in other words,  
an interface should be represented  
by an averaged medium  
consistent with boundary conditions on the interface

Have a quick look at history:

strange (though not unusual) but illuminating and instructive



recall:

an interface should be represented  
by an averaged medium  
consistent with boundary conditions on the interface

Was this the basis for developing heterogeneous FD schemes ?

No, it was not.

Is this now the basis for developing heterogeneous FD schemes ?

Very rarely.

recall:

an interface should be represented  
by an averaged medium  
consistent with boundary conditions on the interface

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Is this now the basis for developing heterogeneous FD schemes ?

Very rarely.

Why?

The explanation would be rather critical  
and related to

persisting unawareness, overlooking and need to sell own schemes and results  
( *sorry* )

All this is supported by arguments on uncertainties of different kinds.

The historical development  
was strongly influenced by the fact that

since 60's until 1984-1988 (introduction of a staggered grid)  
developers FD-approximated  
the 2<sup>nd</sup>-order equation of motion in displacement  
which certainly was not an easy task

( the only reasonable help  
came from Tikhonov and Samarskii  
who used a mathematical trick  
to avoid spatial differentiation of elastic moduli )

After introduction  
of the 1<sup>st</sup>-order velocity-stress formulation on a staggered grid

(which removed the problem of the 2<sup>nd</sup>-spatial derivatives of elastic moduli)

most developers assumed (and explicitly wrote)  
that

the problem of implementing any heterogeneity  
is solved implicitly

—

which obviously is not true

In parallel,

there had been another historical development

since 60's ...

Backus (1962) found out  
how to replace a stack of finely layered medium  
by an averaged medium  
consistent with boundary conditions  
at the interfaces between layers.

Schoenberg & Muir (1989) extended the Backus approach  
to arbitrary anisotropic layers.

Neither Backus nor Schoenberg & Muir  
mentioned a relation to FD modelling.

Explicit essential reference to FD modelling  
was done by Muir et al. (1992).

Nevertheless,  
until the article by  
Moczko, Kristek, Vavryčuk, Archuleta and Halada (2002),  
the concept of an averaged medium  
consistent with the interface boundary conditions  
did not impact  
the heterogeneous FD schemes.

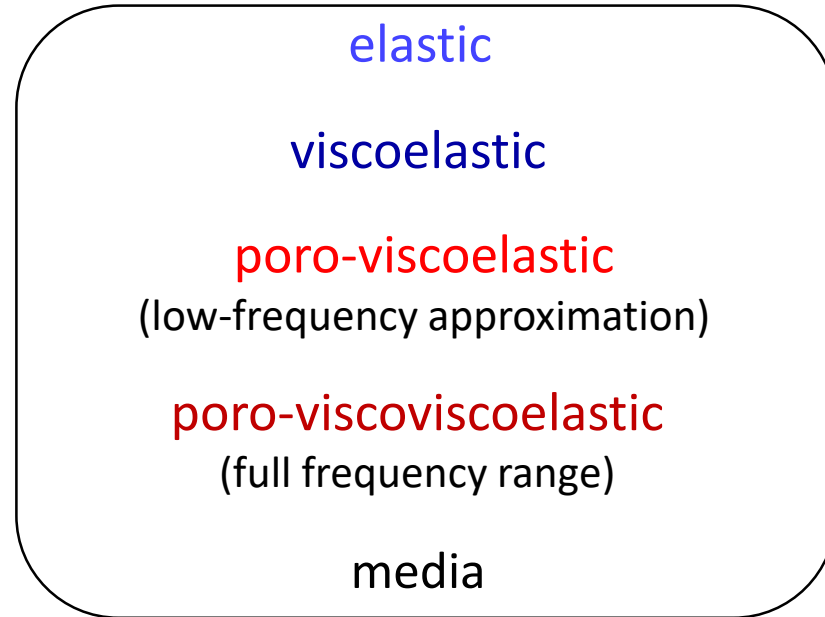
Even then  
many modellers (keen to sell own schemes?)  
have been somehow “overlooking”  
the evident and necessary progress.

This concludes  
the short history of two non-communicating developments 😊

[Back to the present](#)



We have developed  
unified discrete representations  
of a strong material heterogeneity for



The representations have capability of  
a sub-cell resolution  
and thus allow for  
an arbitrary shape and position of an interface in a grid

## Our principles

of finding an averaged medium  
consistent with boundary conditions on an interface

### 1

The stiffness matrix of the averaged medium  
has to have the same structure  
(the same number of nonzero elements)  
as the stiffness matrix of a smooth medium has  
( except that the number of independent elements may be different)

#### consequence:

the number of algebraic operations  
for updating stress  
is the same as for the smooth medium

## Our principles

of finding an averaged medium  
consistent with boundary condition on an interface

### 2

If a grid cell contains a planar interface  
between two homogeneous materials  
perpendicular to a coordinate axis,  
the stiffness matrix of the averaged medium in the cell  
will correspond to the exactly averaged medium  
accounting correctly  
for the position of the interface in the grid cell.

# Example

2D P-SV problem

poroelastic medium (low-frequency approximation)

boundary conditions

continuity of the

traction vector

fluid pressure

solid displacement vector

normal component

of the relative fluid displacement vector

# constitutive law

smooth

$$\begin{bmatrix} \sigma_{xx} \\ \sigma_{zz} \\ \sigma_{xz} \\ -p \end{bmatrix} = \begin{bmatrix} \Lambda + \alpha^2 M & \lambda + \alpha^2 M & 0 & \alpha M \\ \lambda + \alpha^2 M & \Lambda + \alpha^2 M & 0 & \alpha M \\ 0 & 0 & 2\mu & 0 \\ \alpha M & \alpha M & 0 & M \end{bmatrix} \begin{bmatrix} \varepsilon_{xx} \\ \varepsilon_{zz} \\ \varepsilon_{xz} \\ \varepsilon_w \end{bmatrix}$$

averaged

$$\begin{bmatrix} \sigma_{xx} \\ \sigma_{zz} \\ \sigma_{xz} \\ -p \end{bmatrix} = \begin{bmatrix} XX + \frac{XP \ XP}{\Psi} & XZ + \frac{XP \ ZP}{\Psi} & 0 & \frac{XP}{\Psi} \\ XZ + \frac{XP \ ZP}{\Psi} & ZZ + \frac{ZP \ ZP}{\Psi} & 0 & \frac{ZP}{\Psi} \\ 0 & 0 & 2\langle \mu \rangle^{Hx} & 0 \\ \frac{XP}{\Psi} & \frac{ZP}{\Psi} & 0 & \frac{1}{\Psi} \end{bmatrix} \begin{bmatrix} \varepsilon_{xx} \\ \varepsilon_{zz} \\ \varepsilon_{xz} \\ \varepsilon_w \end{bmatrix}$$

where, for example,

$$XX = \left\langle \left\langle \Lambda - \frac{\lambda^2}{\Lambda} \right\rangle^z + \left( \left\langle \frac{\lambda}{\Lambda} \right\rangle^z \right)^2 \langle \Lambda \rangle^{Hx} \right\rangle^{Hx} \quad XZ = \left\langle \frac{\lambda}{\Lambda} \right\rangle^{xz} \langle \Lambda \rangle^{Hxz}$$

equation of motion

smooth

$$\begin{bmatrix} \dot{v}_x \\ -\dot{q}_x \\ \dot{v}_z \\ -\dot{q}_z \end{bmatrix} = \frac{1}{\begin{pmatrix} \rho & \rho_f \\ \rho_f & m \end{pmatrix}} \begin{bmatrix} \frac{1}{\rho} & \frac{1}{m} & \frac{b}{m} & 0 & 0 & 0 \\ \frac{1}{m} & \frac{\rho}{\rho_f} \frac{1}{m} & \frac{\rho}{\rho_f} \frac{b}{m} & 0 & 0 & 0 \\ 0 & 0 & 0 & \frac{1}{\rho} & \frac{1}{m} & \frac{b}{m} \\ 0 & 0 & 0 & \frac{1}{m} & \frac{\rho}{\rho_f} \frac{1}{m} & \frac{\rho}{\rho_f} \frac{b}{m} \end{bmatrix} \begin{bmatrix} \sigma_{xx',x} + \sigma_{xz',z} \\ p_{,x} \\ q_x \\ \sigma_{zz',z} + \sigma_{xz',x} \\ p_{,z} \\ q_z \end{bmatrix}$$

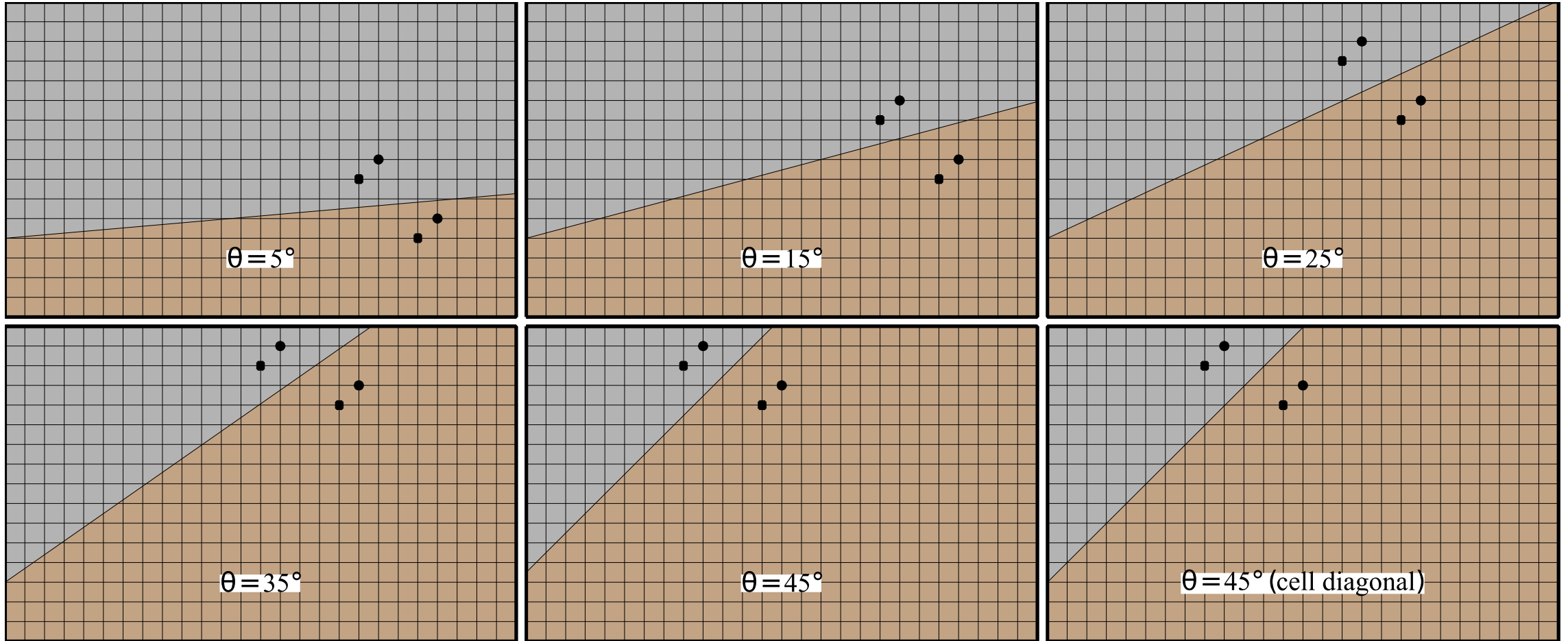
averaged

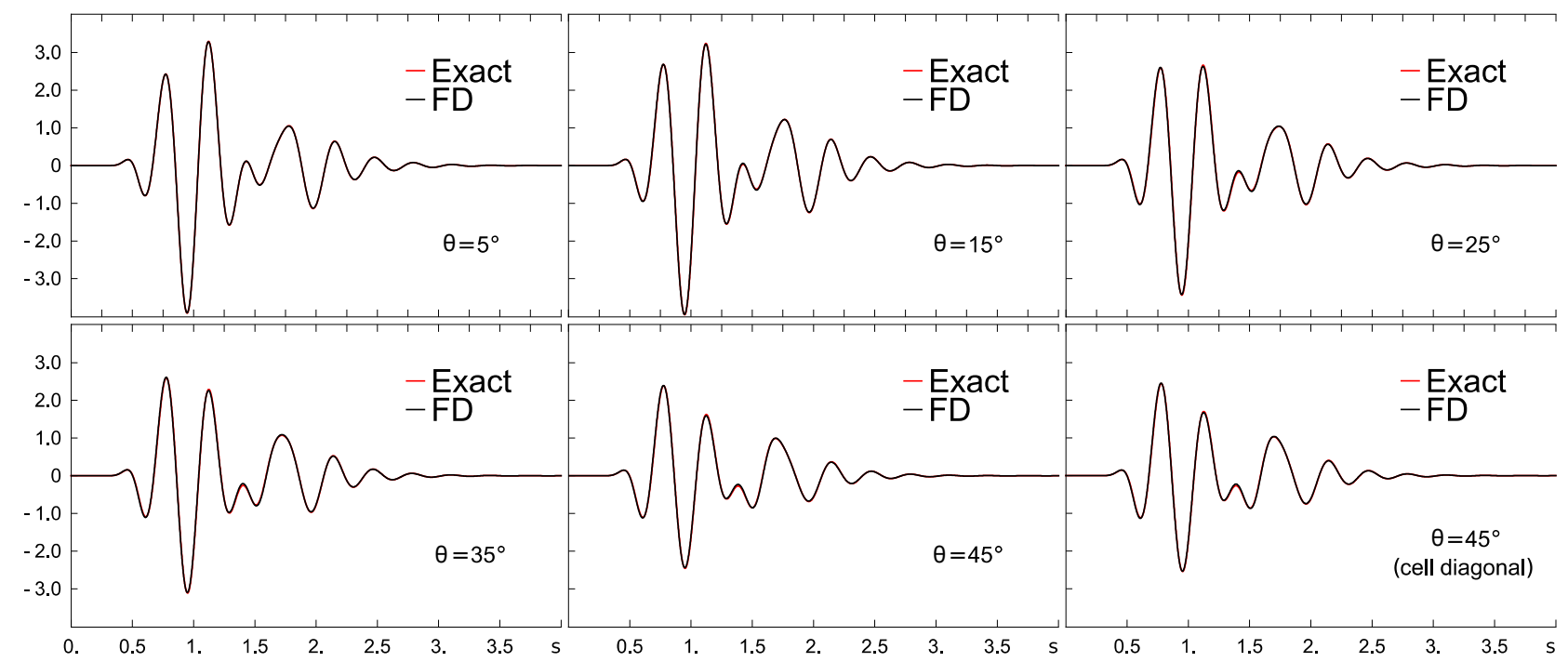
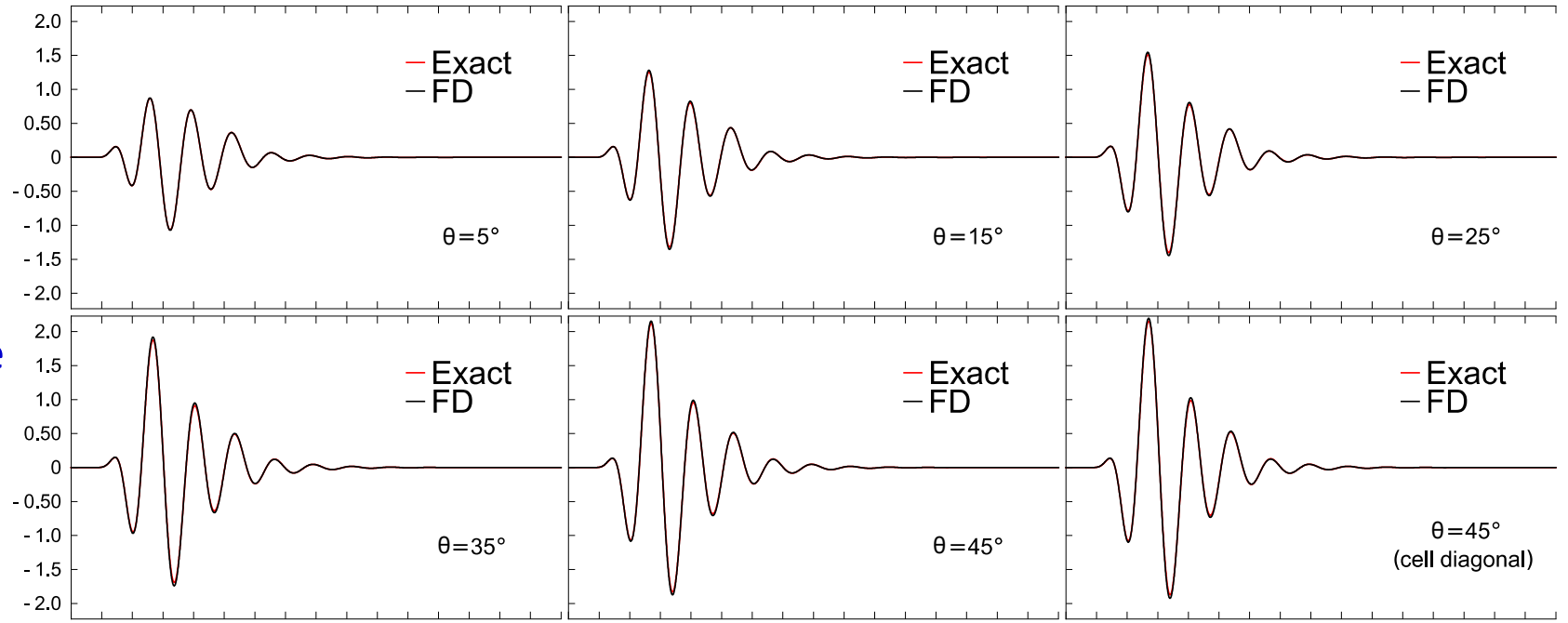
$$\begin{bmatrix} \dot{v}_x \\ -\dot{q}_x \\ \dot{v}_z \\ -\dot{q}_z \end{bmatrix} = \begin{bmatrix} \frac{\langle F^x \rangle^z}{\langle S^x \rangle^z} & \frac{\langle G^x \rangle^z}{\langle S^x \rangle^z} & \frac{\langle H^x \rangle^z}{\langle S^x \rangle^z} & 0 & 0 & 0 \\ \frac{\langle R^x \rangle^z \langle G^x \rangle^z}{\langle S^x \rangle^z} & \frac{\langle P^x \rangle^z \langle G^x \rangle^z}{\langle S^x \rangle^z} & \frac{\langle P^x \rangle^z \langle H^x \rangle^z}{\langle S^x \rangle^z} & 0 & 0 & 0 \\ 0 & 0 & 0 & \frac{\langle F^z \rangle^x}{\langle S^z \rangle^x} & \frac{\langle G^z \rangle^x}{\langle S^z \rangle^x} & \frac{\langle H^z \rangle^x}{\langle S^z \rangle^x} \\ 0 & 0 & 0 & \frac{\langle R^z \rangle^x \langle G^z \rangle^x}{\langle S^z \rangle^x} & \frac{\langle P^z \rangle^x \langle G^z \rangle^x}{\langle S^z \rangle^x} & \frac{\langle P^z \rangle^x \langle H^z \rangle^x}{\langle S^z \rangle^x} \end{bmatrix} \begin{bmatrix} \sigma_{xx',x} + \sigma_{xz',z} \\ p_{,x} \\ q_x \\ \sigma_{zz',z} + \sigma_{xz',x} \\ p_{,z} \\ q_z \end{bmatrix}$$

where, for example,

$$F^\xi = \frac{1}{\langle \rho_f \rangle^\xi} \quad H^\xi = \frac{\langle b \rangle^\xi}{\langle m \rangle^\xi}$$

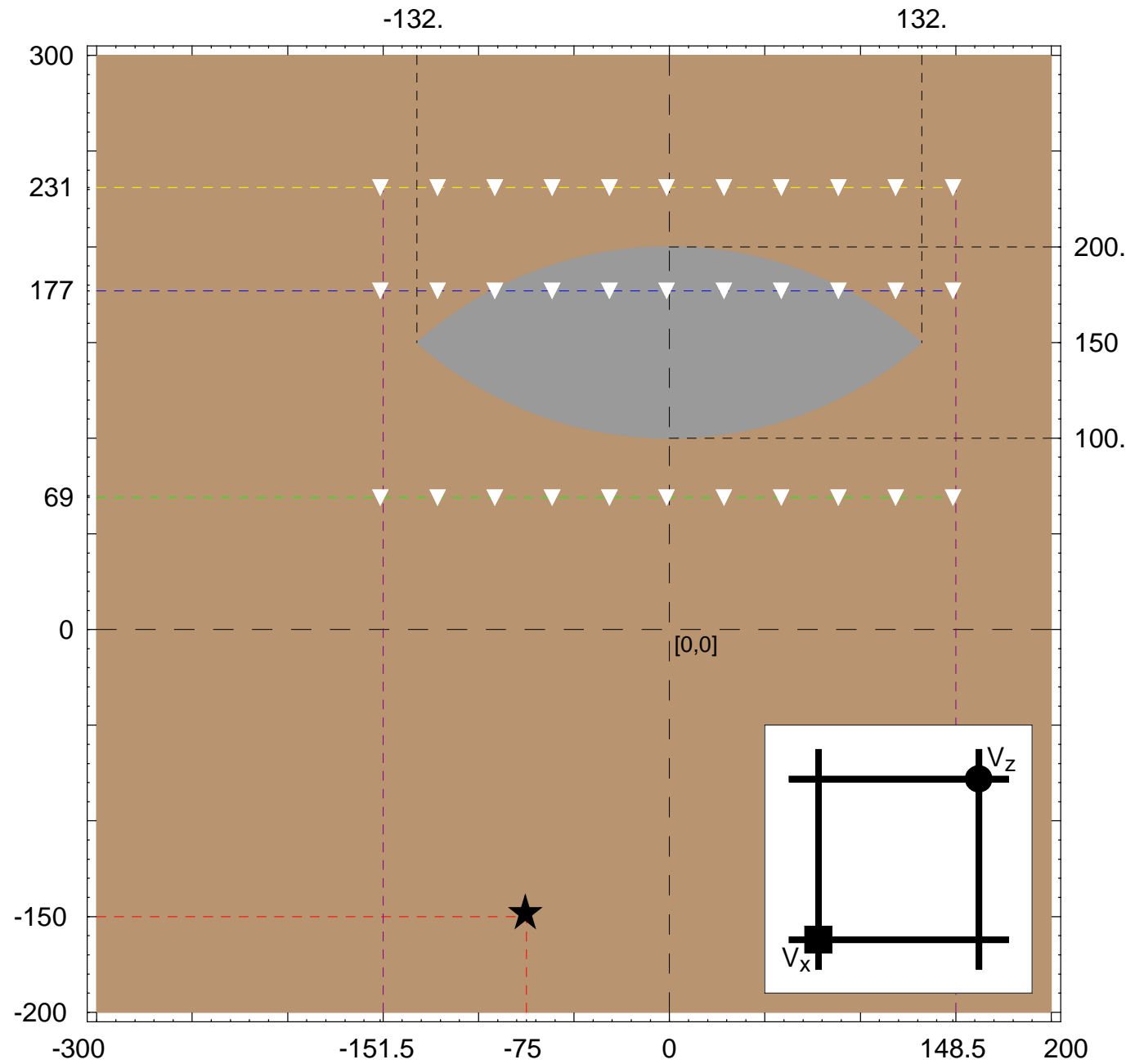
poroelastic halfspaces with an oblique planar interface





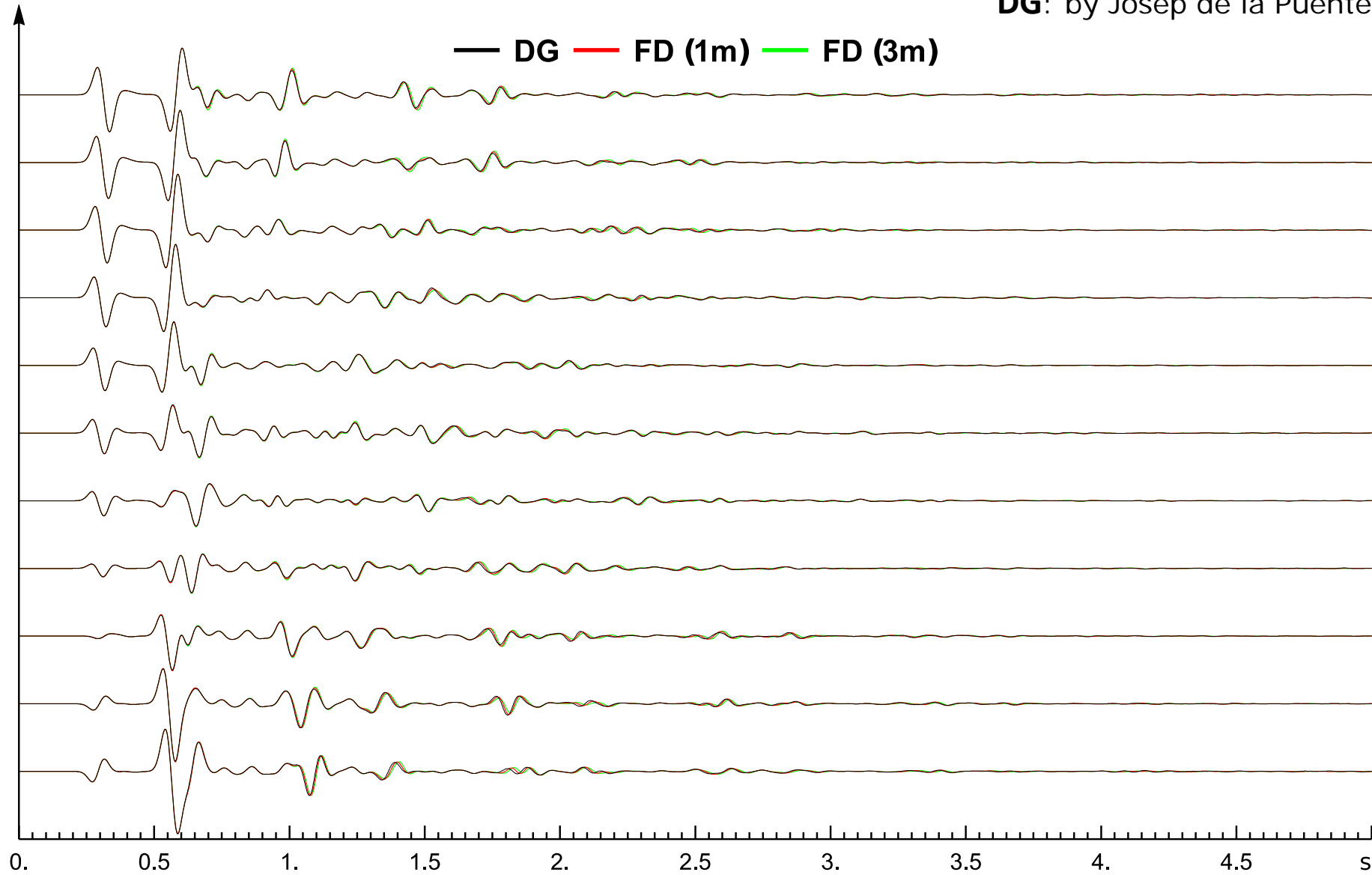


poroelastic lens  
in a poroelastic halfspace



$v_x$  middle-row receivers

DG: by Josep de la Puente

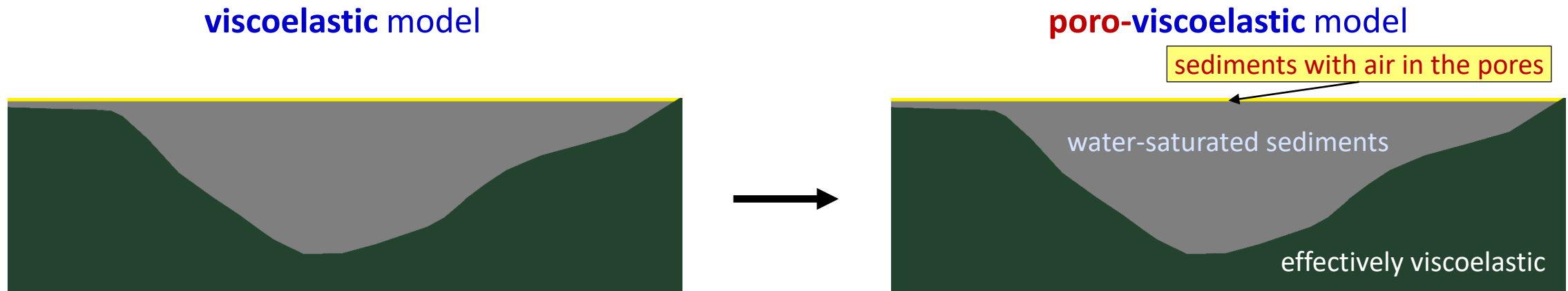


We have performed detailed comparisons  
against SEM (Emmanuel Chaljub, Florent De Martin)  
for  
stringent **viscoelastic** models specially designed  
for testing accuracy, sub-cell resolution and computational efficiency  
of our discrete representation

The models included  
sophisticated canonical configurations  
and  
complex models of the Mygdonian basin near Thessaloniki

**The comparisons confirmed**  
accuracy, sub-cell resolution and computational efficiency  
of our discrete representation

Because we have  
an efficient and sufficiently accurate  
representation of heterogeneity of the **poro-viscoelastic** medium  
we can apply this representation  
also to  
models with **viscoelastic** and **poro-viscoelastic** parts



this can be achieved by

- setting appropriate values of porosity and constant permeability
- choosing an appropriate fluid
- setting an appropriate value of the bulk modulus for a solid phase
- calculating solid-phase density from the density of bedrock
- calculating shear and bulk moduli
- setting an appropriate tortuosity

## conclusions

with a sufficiently accurate discrete representation of a material heterogeneity

the most advanced FD schemes

are more efficient

for modelling earthquake ground motion

**in local surface sedimentary structures**

than the spectral-element and discontinuous-Galerkin methods

# Thank you for your attention

Moczo, Kristek, Gális 2014 The Finite-difference Modelling of Earthquake Motions: Waves and Ruptures  
*Cambridge University Press*

Kristek, Moczo, Chaljub, Kristekova 2017 An orthorhombic representation of a heterogeneous medium for the finite-difference modelling of seismic wave propagation  
*Geophys. J. Int.* 208, 1250–126

Moczo, Gregor, Kristek, de la Puente 2019 A discrete representation of material heterogeneity for the finite-difference modelling of seismic wave propagation in a poroelastic medium  
*Geophys. J. Int.* 216, 1072-1099.

Kristek, Moczo, Chaljub, Kristekova 2019 A discrete representation of a heterogeneous viscoelastic medium for the finite-difference modelling of seismic wave propagation  
*Geophys. J. Int.*, in press.

## EXAMPLE: 2D P-SV constitutive law and equations of motion

poroelastic medium (low-frequency approximation)

$$\begin{bmatrix} \sigma_{xx} \\ \sigma_{zz} \\ \sigma_{xz} \\ -p \end{bmatrix} = \begin{bmatrix} \Lambda + \alpha^2 M & \lambda + \alpha^2 M & 0 & \alpha M \\ \lambda + \alpha^2 M & \Lambda + \alpha^2 M & 0 & \alpha M \\ 0 & 0 & 2\mu & 0 \\ \alpha M & \alpha M & 0 & M \end{bmatrix} \begin{bmatrix} \varepsilon_{xx} \\ \varepsilon_{zz} \\ \varepsilon_{xz} \\ \varepsilon_w \end{bmatrix}$$

$$\begin{aligned} \rho \dot{v}_x &= \sigma_{xx,x} + \sigma_{xz,z} - \rho_f \dot{q}_x \\ \rho \dot{v}_z &= \sigma_{xz,x} + \sigma_{zz,z} - \rho_f \dot{q}_z \\ m \dot{q}_x &= -p_{,x} - \rho_f \dot{v}_x - b q_x \\ m \dot{q}_z &= -p_{,z} - \rho_f \dot{v}_z - b q_z \end{aligned}$$

$\sigma_{xx}, \sigma_{zz}, \sigma_{xz}$  total stress-tensor components

$p$  fluid pressure

$\varepsilon_{xx}, \varepsilon_{zz}, \varepsilon_{xz}$  solid matrix strain-tensor components

$w_x, w_z$ ;  $\varepsilon_w \equiv w_{x,x} + w_{z,z}$  components of displacement of the fluid relative to the solid frame

$\Lambda \equiv \lambda + 2\mu$  Lamé elastic coefficients of the solid matrix

$\alpha$  poroelastic coefficient of effective stress

$M$  coupling modulus between the solid and fluid

$v_x, v_z$  solid particle velocities

$q_x, q_z$  fluid particle velocities relative to the solid

$\rho, \rho_f$  total and fluid densities

$m$  mass coupling coefficient

$b$  resistive friction

boundary conditions at an interface between  
poroelastic media

continuity of the

$$\sigma_{ij}^+ n_j = \sigma_{ij}^- n_j \quad \text{traction vector}$$

$$p^+ = p^- \quad \text{fluid pressure}$$

$$u_i^+ = u_i^- \quad \text{solid displacement vector}$$

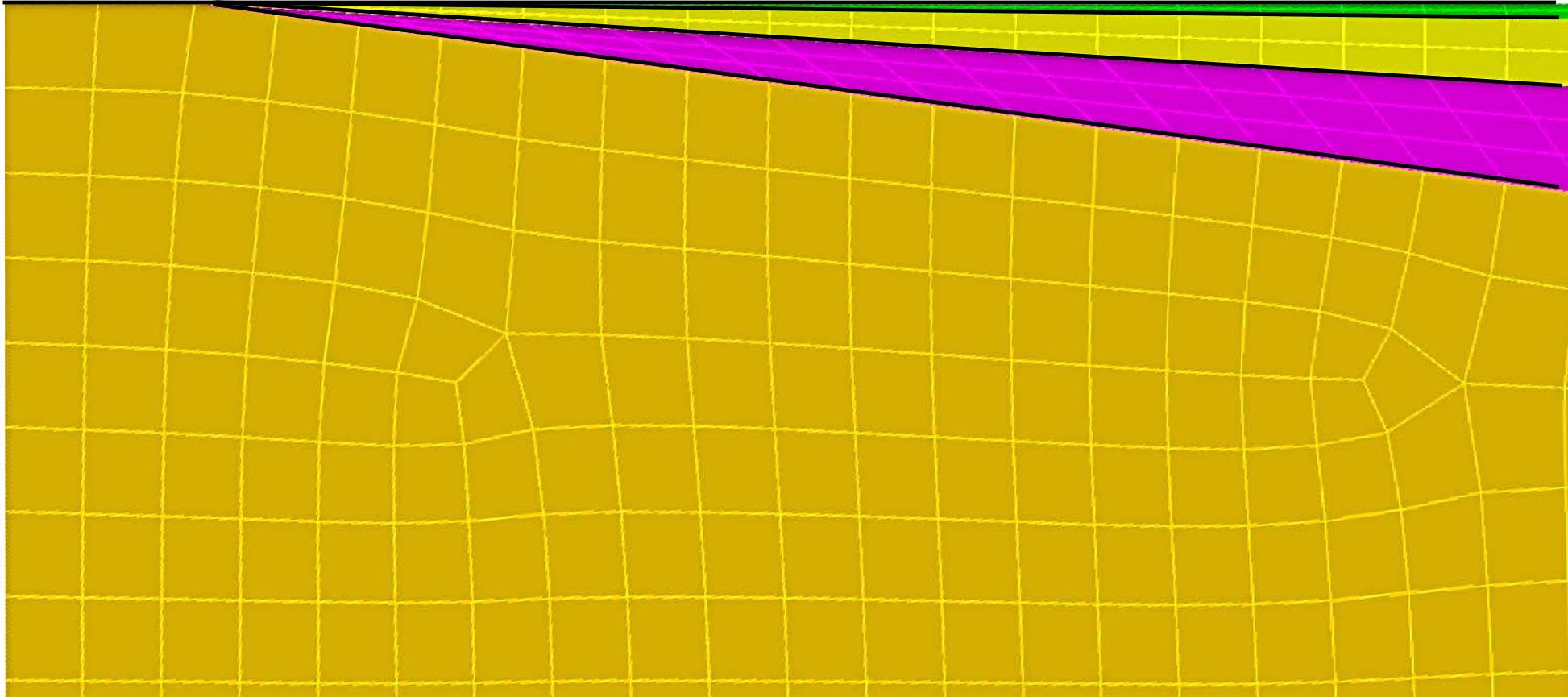
$$w_i^+ n_i = w_i^- n_i \quad \text{normal component} \\ \text{of the relative fluid displacement vector}$$



# interfaces in a sedimentary wedge



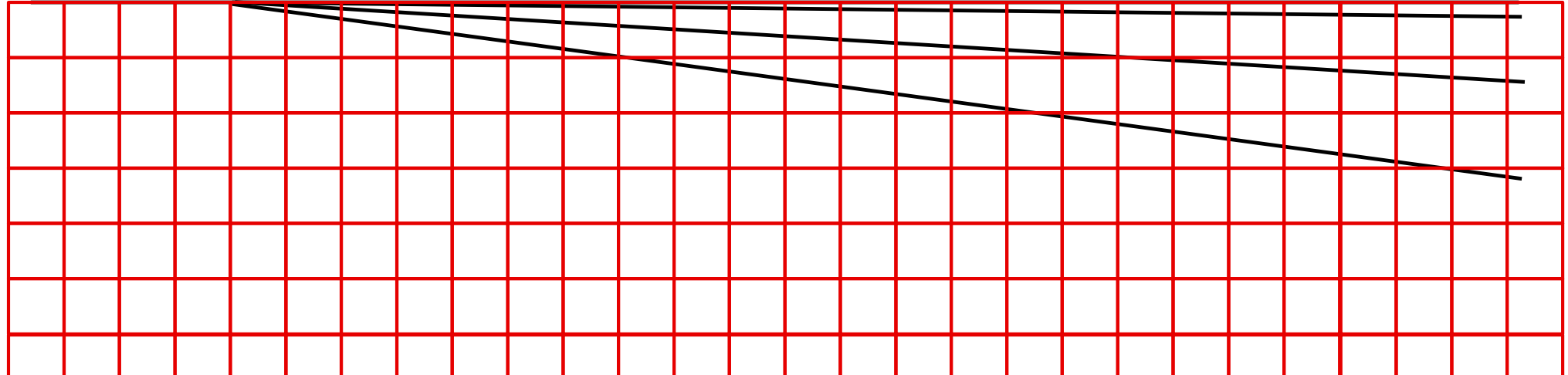
# interfaces in a sedimentary wedge



# interfaces in a sedimentary wedge



# interfaces in a sedimentary wedge



SPEM minimum node-to-node distance: 0.1 m  
FDM grid spacing: 5.0 m